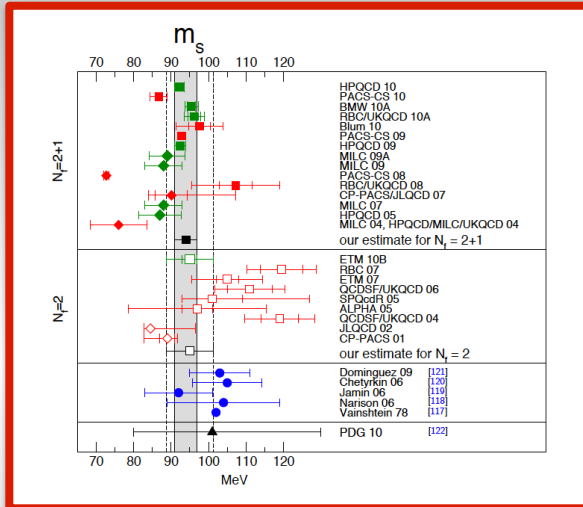


Lattice e flavour nell'era di SuperB

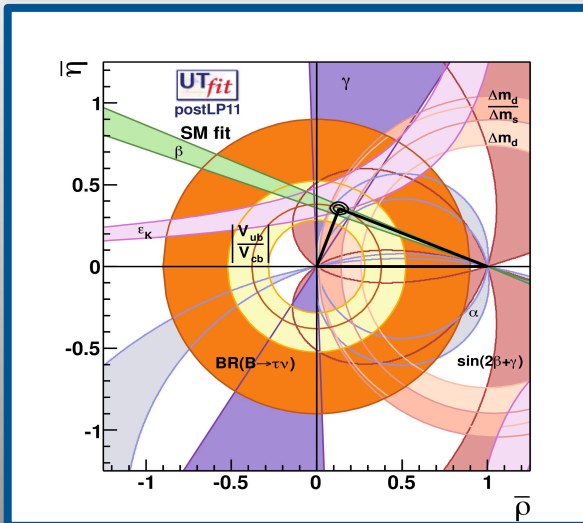
Vittorio
Lubicz



CORTONA 2012

Convegno Informale di Fisica Teorica

30 Maggio - 1 Giugno 2012



GRAZIE

AL COMITATO
ORGANIZZATORE:

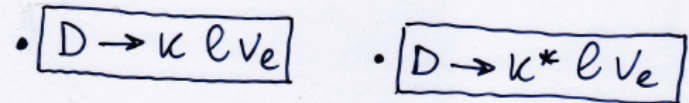
Marisa Bonini, Raffaella
Burioni, Giovanni Cicuta,
Francesco DiRenzo, Luca
Griguolo, Enrico Onofri,
Ettore Vicari

CORTONA
1990

FATTORI DI FORMA E COSTANTI DI DECADIMENTO DEI MESONI D E B NELLA QCD SUL RETICOLO

Collaborazione ELC
(C.R. Allton, M. Guisafulli, V.L.,
L. Maiani, G. Montarilli, C.T. Sachrajda)

1) DECADIMENTI SEMILEPTONICI DEI MESONI D



2) MISURA DI f_B

[IMPORTANTE PER IL $B^0 - \bar{B}^0$ MIXING
 $\langle B^0 | O^{(\Delta B=2)} | \bar{B}^0 \rangle = \frac{8}{3} f_B^2 m_B^2 B, B \sim 1$]

Lattice e flavour nell'era di SuperB

CORTONA 2012

THE PALAZZONE
IN CORTONA



SOMMARIO

- Motivazioni per la fisica del flavour e ruolo della QCD sul reticolo
- Le masse dei quark u, d, s
- La matrice CKM e il test di unitarietà della prima riga: $V_{us}, f_K/f_\pi, f_+(0)$
- Gli effetti di isospin breaking
- L'analisi del triangolo unitario: il passato, il presente e il futuro (SuperB)

The Standard Model does not:

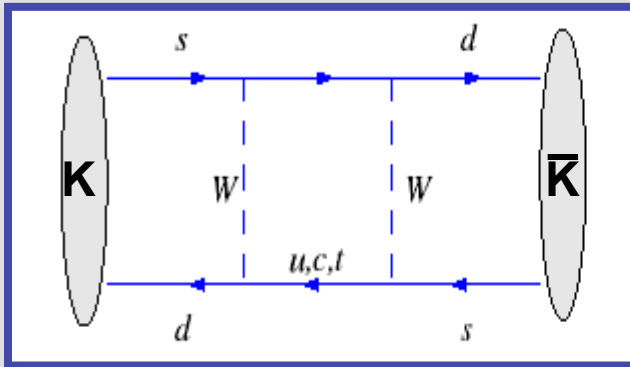
- include gravity ($M_{\text{Planck}} = (\hbar c/G_N)^{1/2} \approx 10^{19} \text{ GeV}$)
- explain the origin of flavour, i.e. masses, mixing and CPV ($M_{\text{Flav}} = ??$)
- give a natural explanation of the hierarchy problem ($M_{\text{Weak}} \ll M_{\text{Planck}}$)
- provide (exact) gauge coupling unification ($M_{\text{GUT}} \approx 10^{15}-10^{16} \text{ GeV}$)
- explain the smallness of neutrino masses ($m_\nu \approx (\lambda v)^2/M$, $M \approx 10^{15} \text{ GeV}$)
- produce the observed matter-antimatter asymmetry
- provide a viable dark matter candidate
- explain "dark" (vacuum) energy



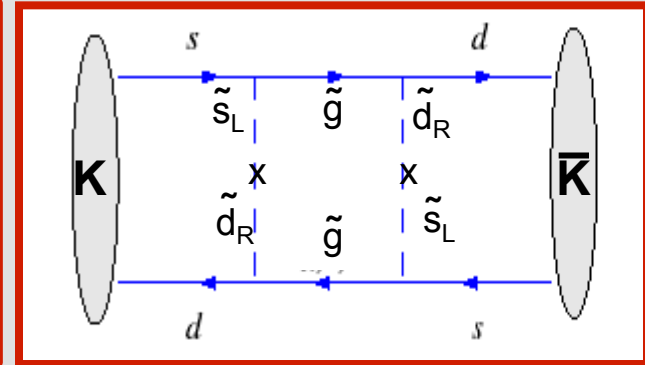
Cosmology

Motivations for flavour physics

Standard Model



New Physics



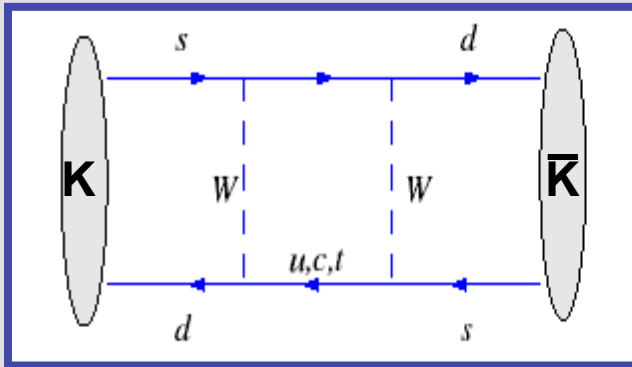
Indirect searches look for new physics (NP) through virtual effects of new particles in loop corrections

- In the SM, FCNC and CP-violating processes occur at the loop level
- In the SM, FV and CPV are governed by the weak interactions and suppressed by mixing angles
- In the SM, quark CPV comes from a single source (neglecting θ_{QCD})

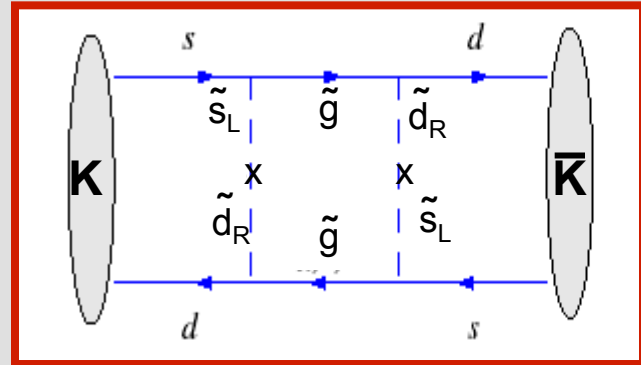
NP does not necessarily share the SM pattern of FV and CPV: **very large NP effects are possible**

Motivations for flavour physics

Standard Model



New Physics



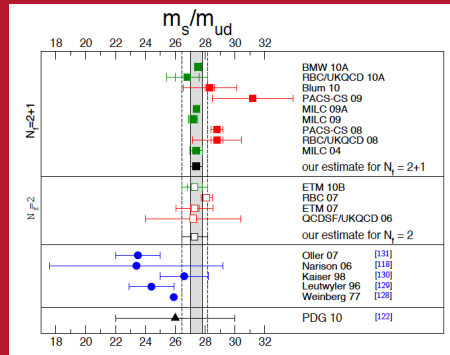
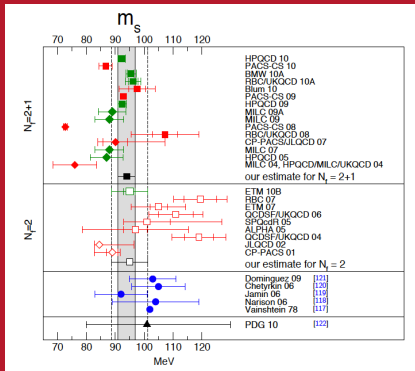
Past (SM) successes in anticipating heavy flavours:

- 1970: **charm** from $K^0 \rightarrow \mu\mu$ (GIM)
- 1973: **3rd generation** from ϵ_K (Kobayashi and Maskawa)
- mid 80s: **heavy top** from semileptonic decays and Δm_B

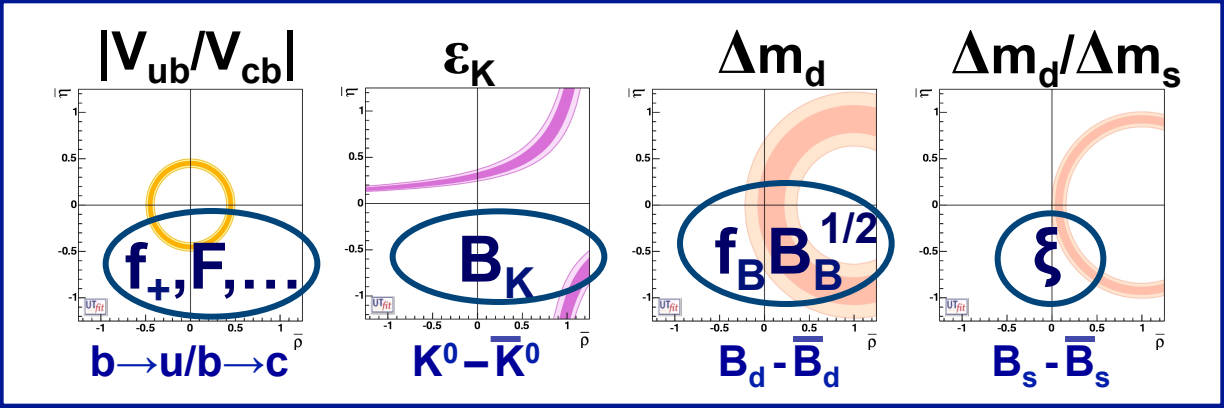
Current and next-generation flavour experiments will improve the experimental precision/sensitivity by one order of magnitude

- **Enough NP-insensitive observables** to pin down the SM contribution with the required accuracy
- **Several NP-sensitive observables** not limited by systematics or theoretical uncertainties

Lattice QCD and flavour physics



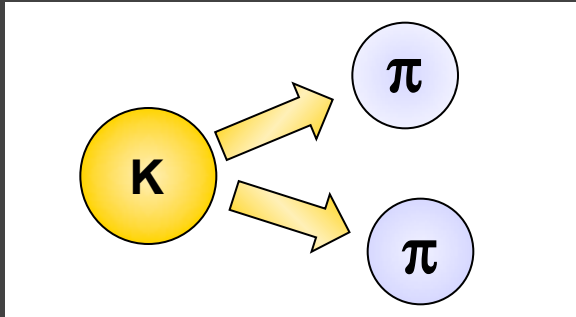
Quark masses



CKM matrix elements

UTA

Beyond SM physics



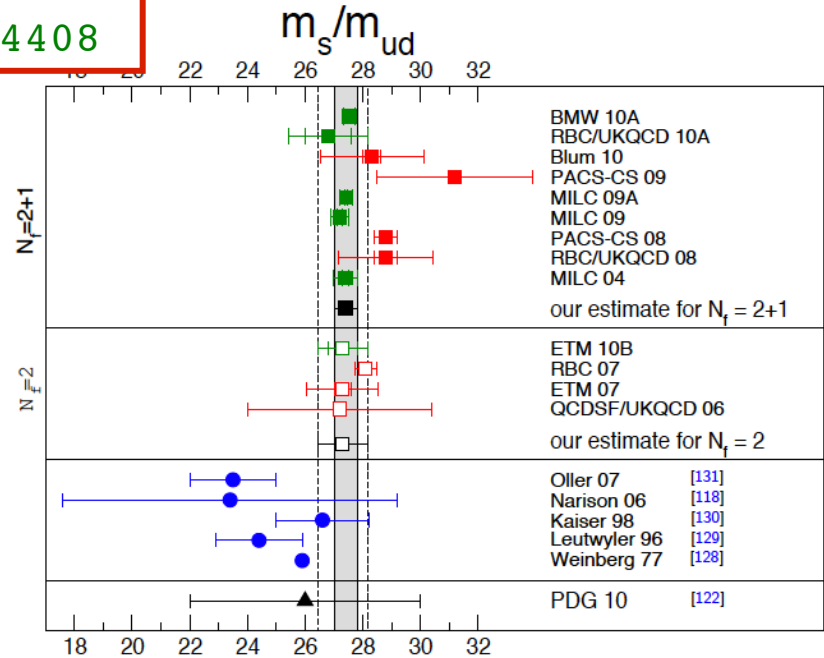
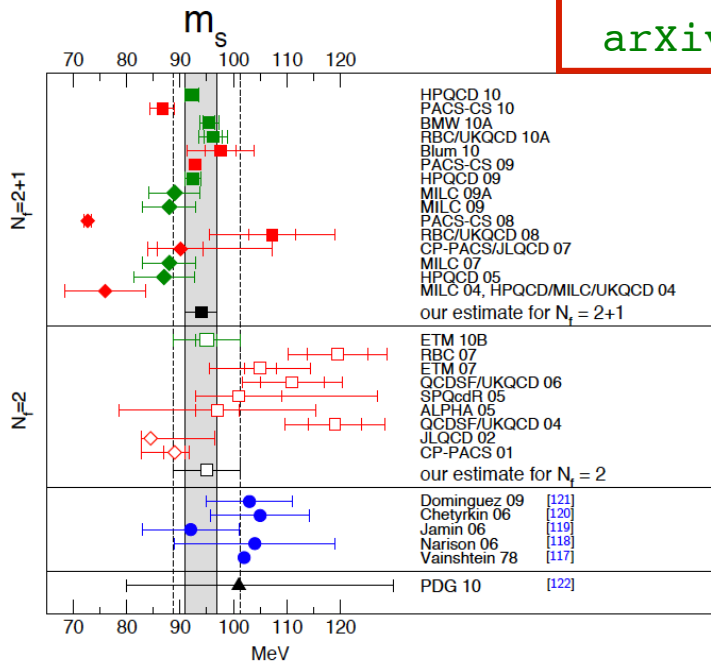
More difficult problems

Not covered in this talk

QUARK MASSES

A benchmark calculation for Lattice QCD

FLAG
arXiv:1011.4408

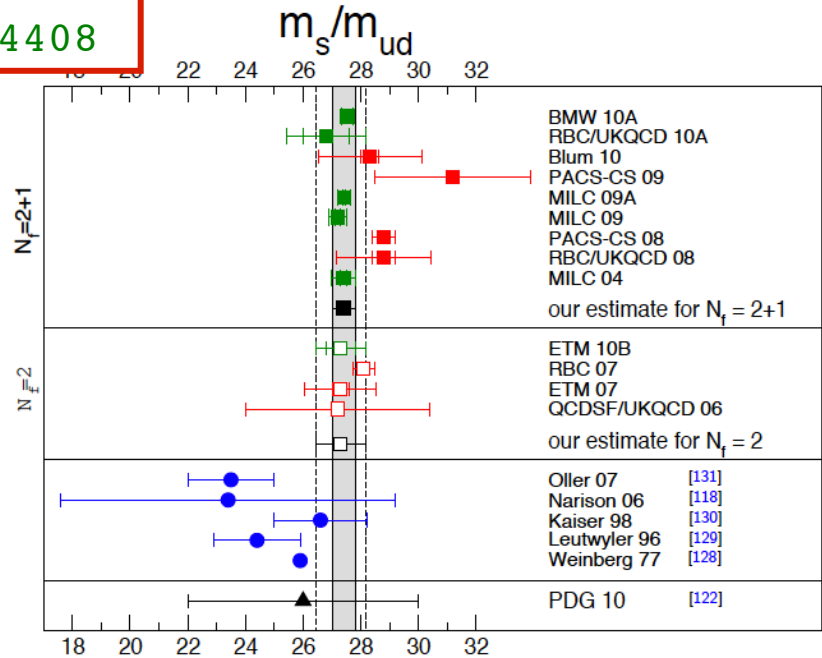
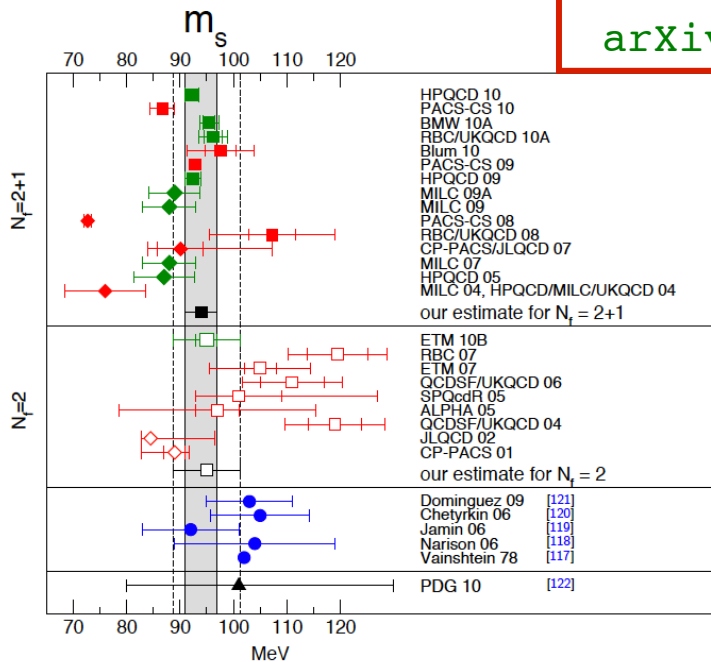


- Many different methods to regularize QCD on the lattice
- Green (Red) : Included (Not included) in the average
- Full (Empty) : $N_f=2+1$ ($N_f=2$) dynamical quarks [$N_f=2+1+1$ in progress]
- Squares (Diamonds) : Non-perturbative (Perturbative) renormalization

QUARK MASSES

A benchmark calculation for Lattice QCD

FLAG
arXiv:1011.4408



$$\bar{m}_s = 94 \pm 3 \text{ MeV} \quad (N_f=2+1)$$

$$\bar{m}_s = 95 \pm 6 \text{ MeV} \quad (N_f=2)$$

$$m_s/m_{ud} = 27.4 \pm 0.4 \quad (N_f=2+1)$$

$$m_s/m_{ud} = 27.3 \pm 0.9 \quad (N_f=2)$$

The accuracy is at the few per cent level

The FLAG colour coding

- Chiral extrapolation:

- ★ $M_{\pi,\min} < 250$ MeV
- $250 \text{ MeV} \leq M_{\pi,\min} \leq 400$ MeV
- $M_{\pi,\min} > 400$ MeV

- Continuum extrapolation:

- ★ 3 or more lattice spacings, at least 2 points below 0.1 fm
- 2 or more lattice spacings, at least 1 point below 0.1 fm
- otherwise

- Finite-volume effects:

- ★ $(M_{\pi}L)_{\min} > 4$ or at least 3 volumes
- $(M_{\pi}L)_{\min} > 3$ and at least 2 volumes
- otherwise

- Renormalization (where applicable):

- ★ non-perturbative
- 2-loop perturbation theory
(with a converging series)
- otherwise

FLAG - Flavour Lattice Averaging Group

An example of
FLAG table

Collaboration	Ref.	publication status	chiral extrapolation	continuum extrapolation	finite volume	renormalization	running	m_{ud}	m_s
PACS-CS 10	[64]	P	★	■	■	★	<i>a</i>	2.78(27)	86.7(2.3)
MILC 10A	[103]	C	●	★	★	●	—	3.19(4)(5)(16)	—
HPQCD 10	[104]	A	●	★	★	★	—	3.39(6)*	92.2(1.3)
BMW 10A, 10B+	[65, 105]	P	★	★	★	★	<i>b</i>	3.469(47)(48)	95.5(1.1)(1.5)
RBC/UKQCD 10A	[106]	P	●	●	★	★	<i>c</i>	3.59(13)(14)(8)	96.2(1.6)(0.2)(2.1)
Blum 10 [†]	[74]	P	●	■	●	★	—	3.44(12)(22)	97.6(2.9)(5.5)
PACS-CS 09	[42]	A	★	■	■	★	<i>a</i>	2.97(28)(3)	92.75(58)(95)
HPQCD 09	[107]	A	●	★	★	★	—	3.40(7)	92.4(1.5)
MILC 09A	[59]	C	●	★	★	●	—	3.25 (1)(7)(16)(0)	89.0(0.2)(1.6)(4.5)(0.1)
MILC 09	[6]	A	●	★	★	●	—	3.2(0)(1)(2)(0)	88(0)(3)(4)(0)
PACS-CS 08	[63]	A	★	■	■	■	—	2.527(47)	72.72(78)
RBC/UKQCD 08	[108]	A	●	■	★	★	—	3.72(16)(33)(18)	107.3(4.4)(9.7)(4.9)
CP-PACS/ JLQCD 07	[109]	A	■	★	★	■	—	3.55(19)(⁺⁵⁶ ₋₂₀)	90.1(4.3)(^{+16.7} _{-4.3})
HPQCD 05	[110]	A	●	●	●	●	—	3.2(0)(2)(2)(0) [‡]	87(0)(4)(4)(0) [‡]
MILC 04, HPQCD/ MILC/UKQCD 04	[77, 111]	A	●	●	●	■	—	2.8(0)(1)(3)(0)	76(0)(3)(7)(0)
ETM 10B	[94]	A	●	★	●	★	<i>a</i>	3.6(1)(2)	95(2)(6)
JLQCD/TWQCD 08A	[95]	A	●	■	■	★	—	4.452(81)(38)(⁺⁰ ₋₂₂₇)	—
RBC 07 [†]	[73]	A	■	■	★	★	—	4.25(23)(26)	119.5(5.6)(7.4)
ETM 07	[49]	A	●	■	●	★	—	3.85(12)(40)	105(3)(9)
QCDSF/ UKQCD 06	[96]	A	■	★	■	★	—	4.08(23)(19)(23)	111(6)(4)(6)
SPQcdR 05	[97]	A	■	●	●	★	—	4.3(4)(^{+1.1} _{-0.0})	101(8)(⁺²⁵ ₋₀)
ALPHA 05	[98]	A	■	●	★	★	<i>b</i>	—	97(4)(18) [§]
QCDSF/ UKQCD 04	[99]	A	■	★	■	★	—	4.7(2)(3)	119(5)(8)
JLQCD 02	[100]	A	■	■	●	■	—	3.223(⁺⁴⁶ ₋₆₉)	84.5(^{+12.0} _{-1.7})
CP-PACS 01	[101]	A	■	■	★	■	—	3.45(10)(⁺¹¹ ₋₁₈)	89(2)(⁺² ₋₆) [*]

FLAG-1

G.Colangelo,
S.Dür, A.Jüttner,
L.Lellouch,
H.Leutwyler,
V.Lubicz,
S.Necco,
C.Sachrajda,
S.Simula,
T.Vladikas,
U.Wenger,
H.Wittig

arXiv:
1011.4408

Uncertainties
are being
carefully
investigated

The CKM matrix

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

A 3x3 unitary matrix which originates from the misalignment in flavour space of the up and down component of the $SU(2)_L$ quark doublet of the Standard Model

In the quark mass eigenstate basis, V_{CKM} appears in the quark charged-current interaction Lagrangian. It is the only source of flavour-changing transitions and CP violation in the SM

$$L^{cc} = \frac{g}{2\sqrt{2}} \sum_{i,j} \bar{u}_i \gamma_\mu (1 - \gamma_5) (V_{CKM})_{ij} d_j W^\mu + h.c.$$

$$\sum_k (V_{CKM}^+)_{ik} (V_{CKM})_{kj} = \delta_{ij}$$

3 diagonal + 6 triangular relations
 \Rightarrow 9 real parameters: 3 angles, 1 phase
+ 5 unphysical phases rotated away by a redefinition of the quark fields

$$\Rightarrow \theta_{12}, \theta_{13}, \theta_{23}, \delta$$

The CKM matrix

$$V_{CKM} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \times \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \times \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \Rightarrow$$

$$V_{CKM} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{13}c_{23} \end{pmatrix}$$

The PDG
"standard"
parametrization

$$\lambda \equiv s_{12}$$

$$A\lambda^2 \equiv s_{23}$$

$$A\lambda^3(\rho - i\eta) \equiv s_{13}e^{-i\delta}$$

Wolfenstein '83
Buras et al., '94

$$V_{CKM} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \bar{\rho} - i\bar{\eta}) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$

$(\bar{\rho}, \bar{\eta})$ is the apex of the UT

$$\bar{\rho} \approx \rho(1 - \lambda^2/2)$$

$$\bar{\eta} \approx \eta(1 - \lambda^2/2)$$

THE CKM 1st ROW UNITARITY TEST

$$|V_{ud}|^2 + |V_{us}|^2 + \cancel{|V_{ub}|^2} = 1$$

The most stringent unitarity test

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

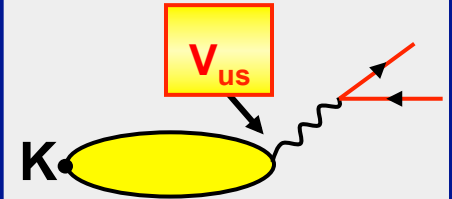
Processes: $K \rightarrow lv$, $K \rightarrow \pi lv$

Theory input: f_K/f_π , $f_+(0)$

V_{us}/V_{ud} from $K\mu 2/\pi\mu 2$ decays

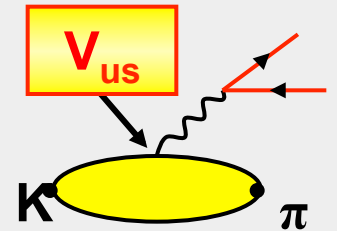
$$\frac{\Gamma(K \rightarrow \mu \bar{\nu}_\mu (\gamma))}{\Gamma(\pi \rightarrow \mu \bar{\nu}_\mu (\gamma))} = \frac{|V_{us}|^2 \left(\frac{f_K}{f_\pi}\right)^2 m_K \left(1 - \frac{m_\mu^2}{m_K^2}\right)}{|V_{ud}|^2 m_\pi \left(1 - \frac{m_\mu^2}{m_\pi^2}\right)} \times 0.9930(35)$$

[Marciano 04]



V_{us} from $Kl3$ decays

$$\Gamma_{K \rightarrow \pi l \nu} = C_K^2 \frac{G_F^2 m_K^5}{192 \pi^2} / S_{EW} [1 + \Delta_{SU(2)} + 2\Delta_{EM}] \times |V_{us}|^2 |f_+^{K\pi}(0)|^2$$



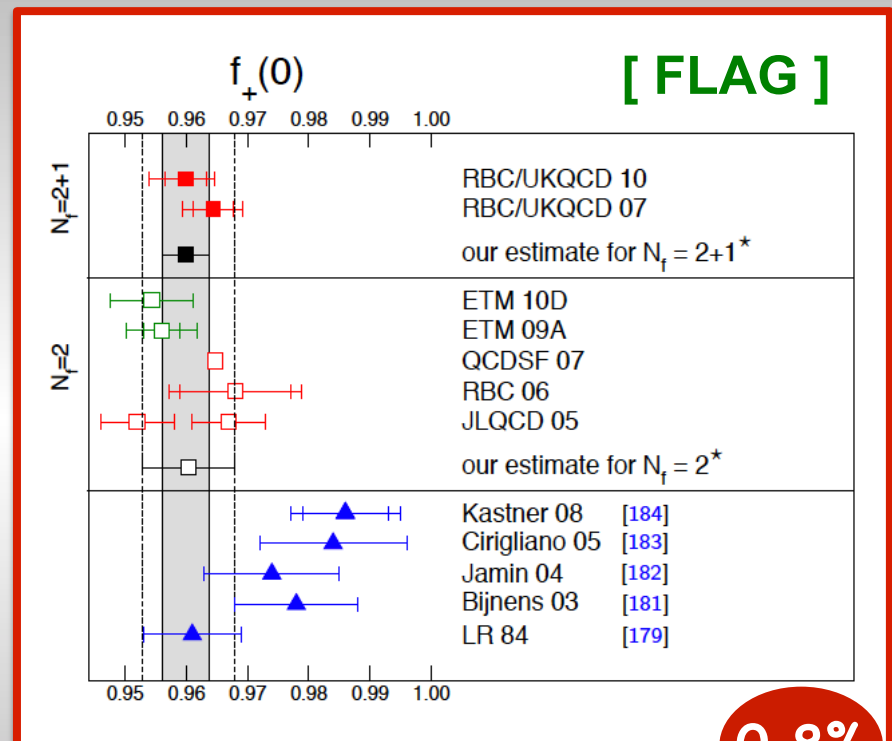
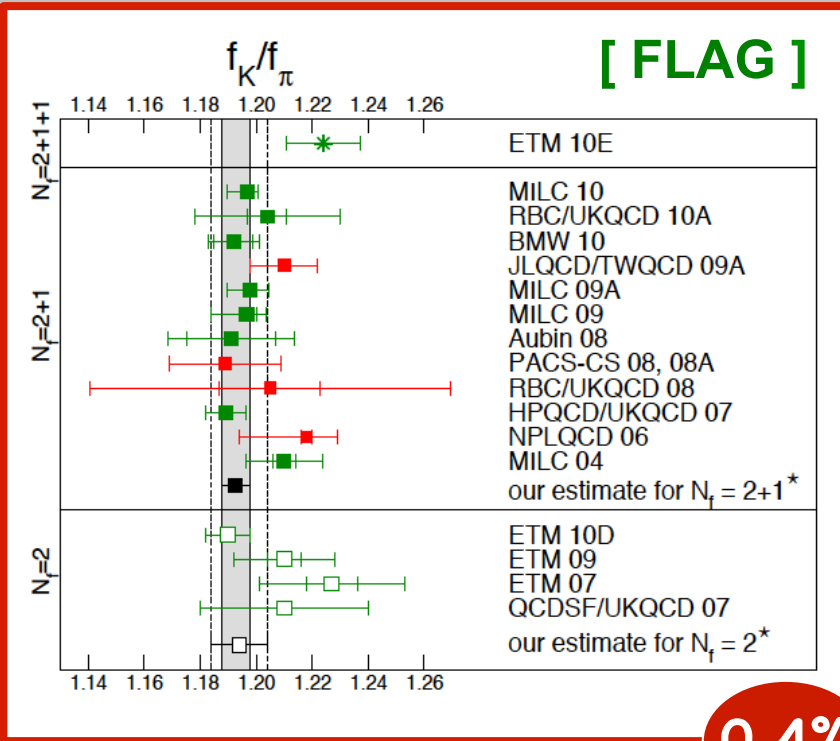
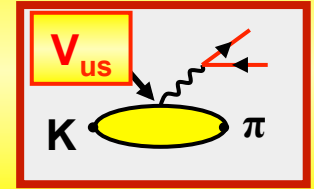
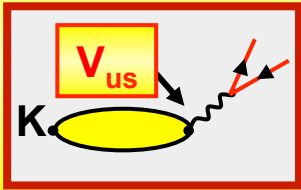
$$\frac{|V_{us}|}{|V_{ud}|} \frac{f_K}{f_\pi} = 0.2758(5)$$

Flavia
net Kaon WG

$$|V_{us}| f_+(0) = 0.2163(5)$$

Lattice results:

f_K/f_π and $f_+(0)$



$f_K/f_\pi = 1.193(5) \quad (N_f=2+1)$

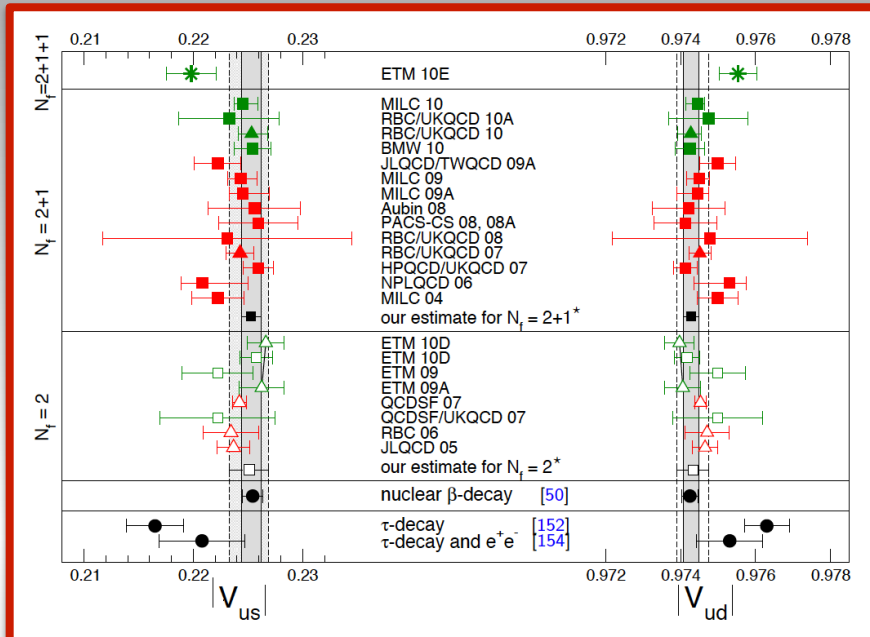
$f_K/f_\pi = 1.210(18) \quad (N_f=2)$

$f_+(0) = 0.956(8) \quad (N_f=2 \text{ \& } 2+1)$

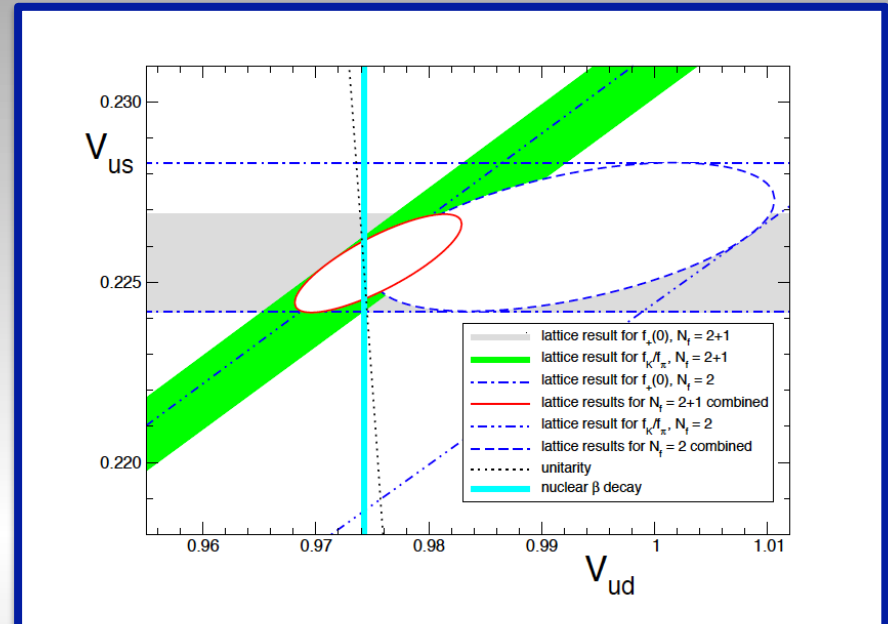
First result with $N_f=2+1+1$ available

Predictions of analytical model tends to be larger than lattice results¹⁶

The 1st row unitarity test



Results from f_K/f_π [squares] and $f_+(0)$ [triangles]



The 1st row unitarity plot [FLAG]

From K_{l2} and K_{l3} decays

$$\bullet |V_{ud}| = 0.9743(2) \quad \bullet |V_{us}| = 0.2254(9)$$

Combining with $|V_{ud}|$ from nuclear β decays:

$$\Delta_{\text{CKM}} = |V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 - 1 = (0 \pm 7) \cdot 10^{-4}$$

ISOSPIN BREAKING EFFECTS

The lattice determinations are usually obtained in the limit of exact ISOSPIN SYMMETRY, i.e. $m_u = m_d$ and $Q_u = Q_d = 0$

E.g. $f_+(0) = 0.956(8)$ 0.8% $f_K/f_\pi = 1.193(5)$ 0.4%

Though small, **isospin breaking effects** are important at the current level of **precision in flavour physics**. Their typical size is:

$Q_u \neq Q_d$: $O(\alpha_{e.m.}) \approx 1/100$

“Electromagnetic”

$m_u \neq m_d$: $O[(m_d - m_u)/\Lambda_{QCD}] \approx 1/100$

“Strong”

down
-1/3

up
+2/3

Isospin breaking effects are also responsible for the stability of matter through the proton-neutron mass difference

The calculation of IB effects on the lattice is challenging 18

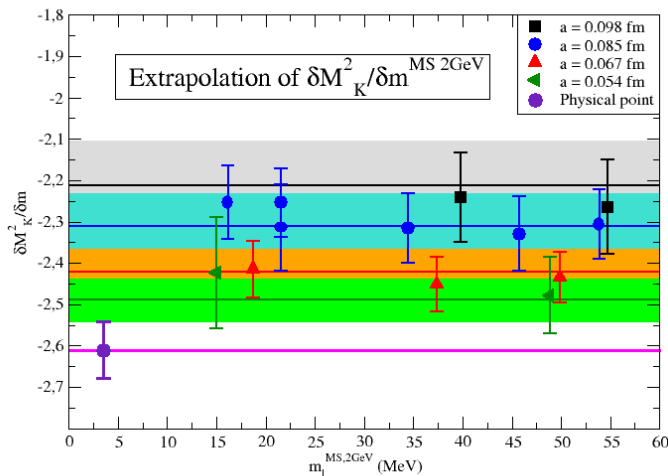
A strategy for Lattice QCD: the (md-mu) expansion

RM123 collaboration, G.M. de Divitiis et al., arXiv:1110.6294

Expand the functional integral in powers of $\delta m = (\mu - m_d)/2$

$$\langle O \rangle \propto \int D\phi O e^{-S_0 + \delta m \hat{S}} \stackrel{1st}{\simeq} \int D\phi O e^{-S_0} (1 + \delta m \hat{S}) = \langle O \rangle_0 + \delta m \langle O \hat{S} \rangle_0$$

The mass difference $m_d - m_u$ is a free parameter of the Lagrangian and one **experimental input** is needed to fix it. A simple choice is the mass splitting between the **neutral and charged kaon**



- We find: $\delta M_K^{QCD} / \delta m^{\overline{MS}, 2\text{GeV}} = 2.64(7)$
- Using as input (from FLAG):

$$(M_{K^0} - M_{K^+})_{QCD} = 2 \delta M_K^{QCD} = 6.0(6)_{QED} \text{ MeV}$$
one obtains:

$$\bar{m}_d - \bar{m}_u = 2.28(6)_{LQCD} (23)_{QED} \text{ MeV}$$

$$\bar{m}_d / \bar{m}_u = 0.51(4)$$

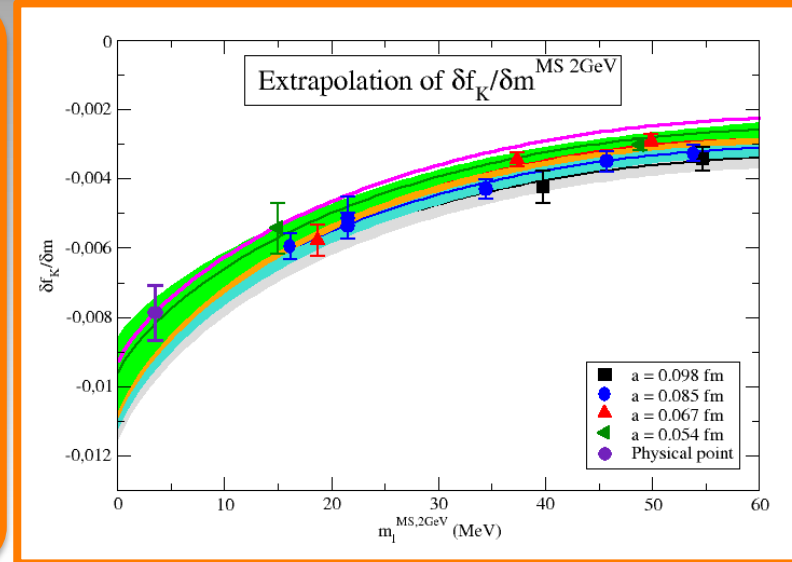
Isospin breaking effects in the ratio f_K/f_π

- We find: $\delta f_K^{QCD} / \delta m^{\overline{MS}, 2\text{GeV}} = -0.47(5)$
- Using the previous result for $m_d - m_u$, we obtain

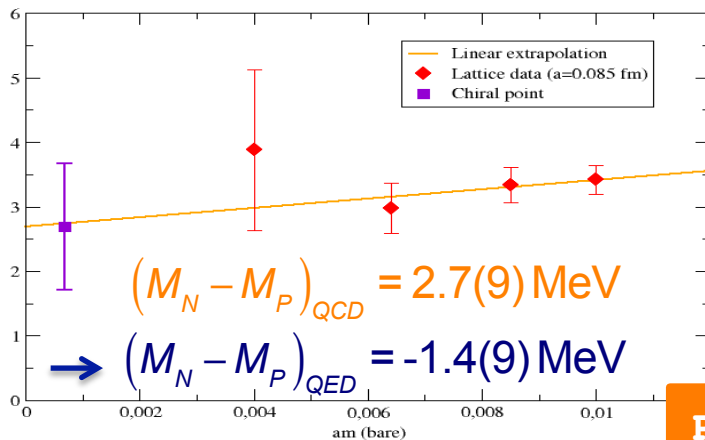
$$\delta(f_K/f_\pi) / (f_K/f_\pi)_{QCD} = -0.0034(3)_{LQCD} (3)_{QED}$$

- It can be compared with the ChPT estimate

$$\delta(f_K/f_\pi) / (f_K/f_\pi)_{QCD} = -0.0022(6)$$

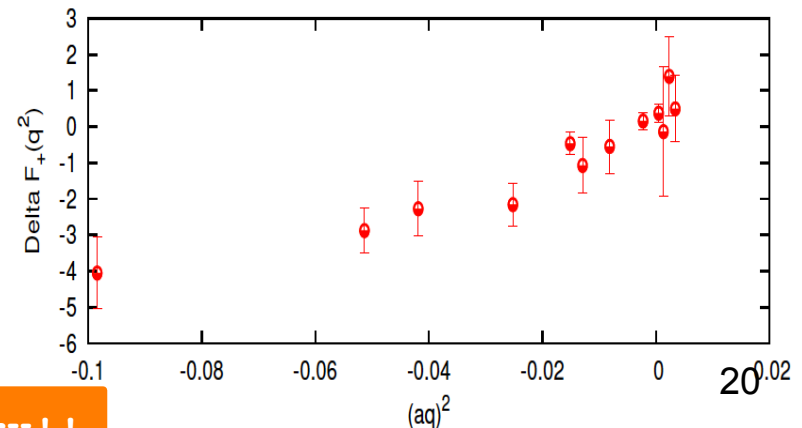


The neutron-proton mass splitting

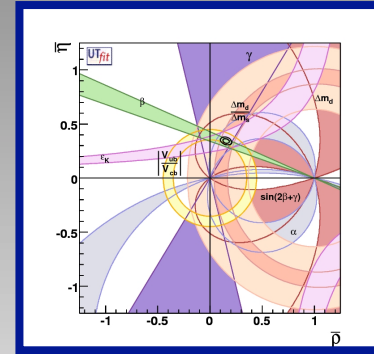


Preliminary!!

The $K \rightarrow \pi l \nu$ semileptonic form factors



THE UNITARITY TRIANGLE ANALYSIS



$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$

$$\left(\begin{array}{ccc} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{array} \right)$$

Processes: $B \rightarrow l \nu$, $B \rightarrow D/\pi l \nu$,

$K - K$, $B_{(s)} - B_{(s)}$

Theory input: f_B , f_{B_s} , $f_+(0)$,

B_K , B_B , ...

The unitarity triangle

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

The relation can be rewritten as:

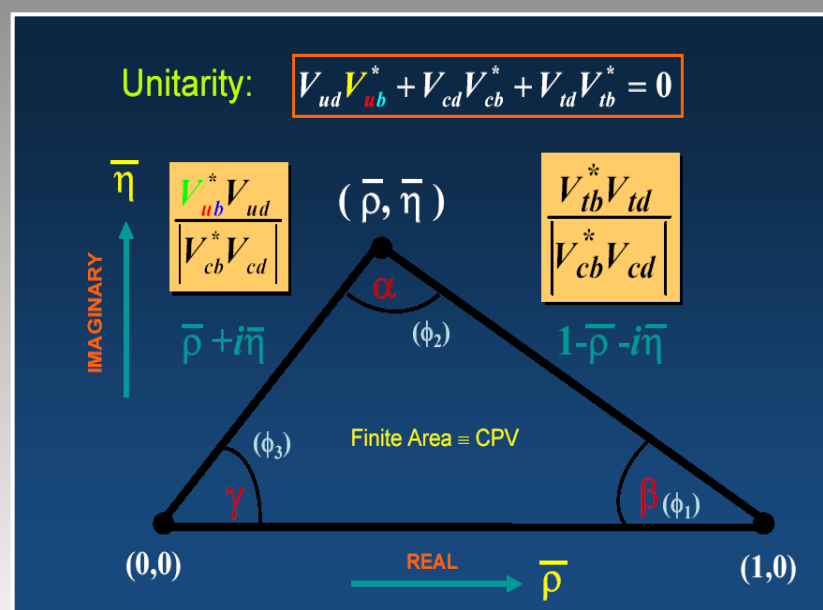
$$R_u e^{i\gamma} + R_t e^{-i\beta} = 1$$

with:

$$R_u = \left| \frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} \right| \quad \gamma = \arg \left(-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} \right)$$

$$R_t = \left| \frac{V_{td}V_{tb}^*}{V_{cd}V_{cb}^*} \right| \quad \beta = \arg \left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*} \right)$$

$$\text{CPV} \sim \text{J} = \text{Im} V_{ij}V_{il}^*V_{kl}V_{kj}^*$$



The apex of the UT defines $(\bar{\rho}, \bar{\eta})$

$$R_u e^{i\gamma} \equiv \bar{\rho} + i\bar{\eta}$$

$$R_t e^{-i\beta} = 1 - \bar{\rho} - i\bar{\eta}$$

$$\begin{aligned} \rho + i\eta &= \sqrt{\frac{1 - A^2 \lambda^4}{1 - \lambda^2}} \frac{\bar{\rho} + i\bar{\eta}}{1 - A^2 \lambda^4 (\bar{\rho} + i\bar{\eta})} \\ &\simeq \left(1 + \frac{\lambda^2}{2} \right) (\bar{\rho} + i\bar{\eta}) + O(\lambda^4) \end{aligned}$$

The



Collaboration



Marco



Enrico



Maurizio



Luca



Cecilia



Vittorio



Guido



Marcella



Vincenzo



Carlo



Denis



Achille



Viola



Adrian



Fabrizio



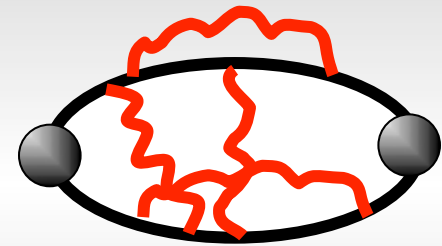
THE UNITARITY TRIANGLE ANALYSIS

① The past
The “quenched” era

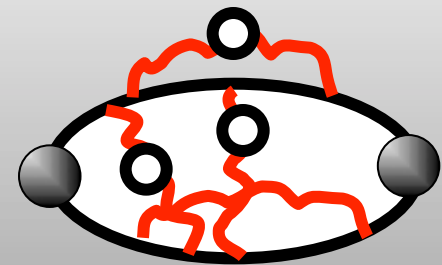
Uncertainties in LQCD before 2006

For many years, uncertainties in lattice calculations have been dominated by the **quenched approximation** (or, more precisely, by the uncertainty on the quenching error)

	f_B [MeV]	$f_{B_s} \sqrt{B_s}$ [MeV]	ξ
J.Flynn Latt' 96	175(25) 14%	----	----
C.Bernard Latt' 00	200(30) 15%	267(46) 17%	1.16(5) 4%
L.Lellouch Ichep' 02	193(27)(10) 15%	276(38) 14%	1.24(4)(6) 6%
Hashimoto Ichep' 04	189(27) 14%	262(35) 13%	1.23(6) 5%
N.Tantalo CKM' 06	223(15)(19) 11%	246(16)(20) 10%	1.21(2)(5) 4%



QUENCHED



UNQUENCHED

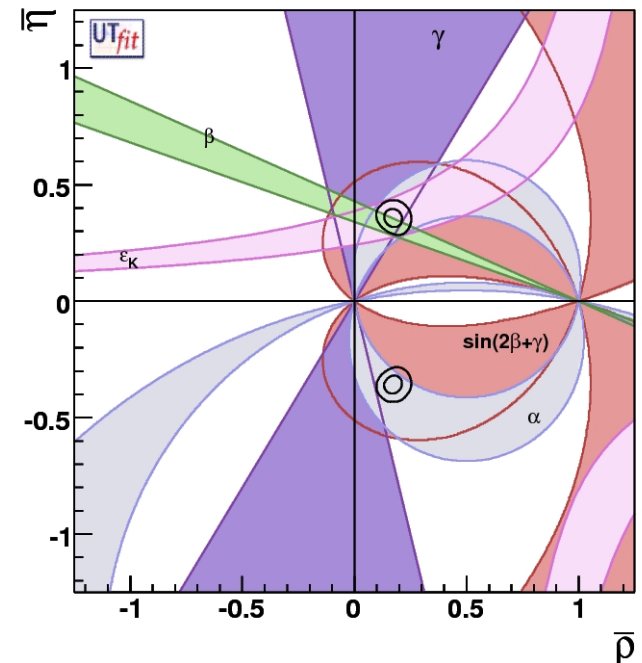
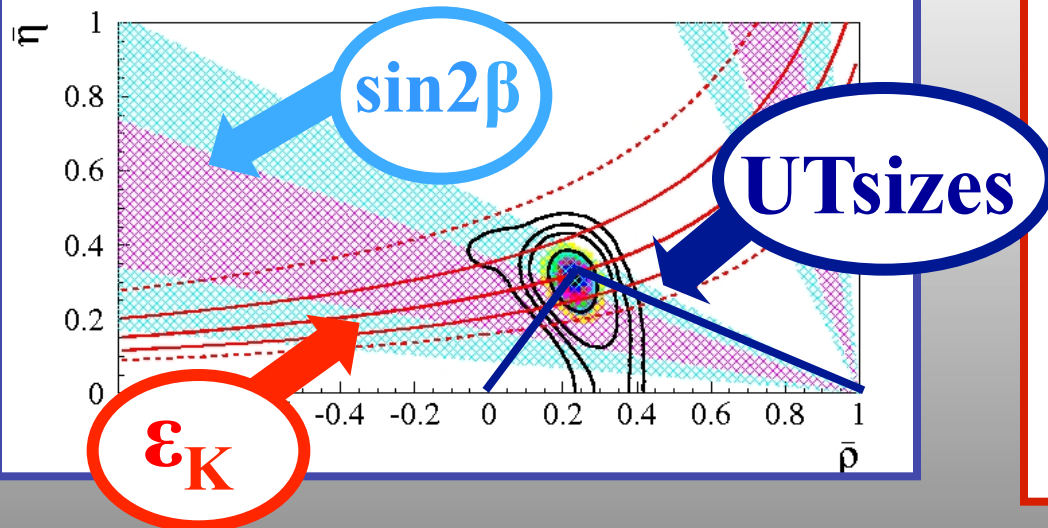
In spite of the relatively large lattice uncertainties, important results for flavour physics have been achieved

CKM PARADIGM OF ~~CP~~

CP-conserving and CP-violating processes determine the same CKM phase

UTfit, today

Ciuchini et al., 2000



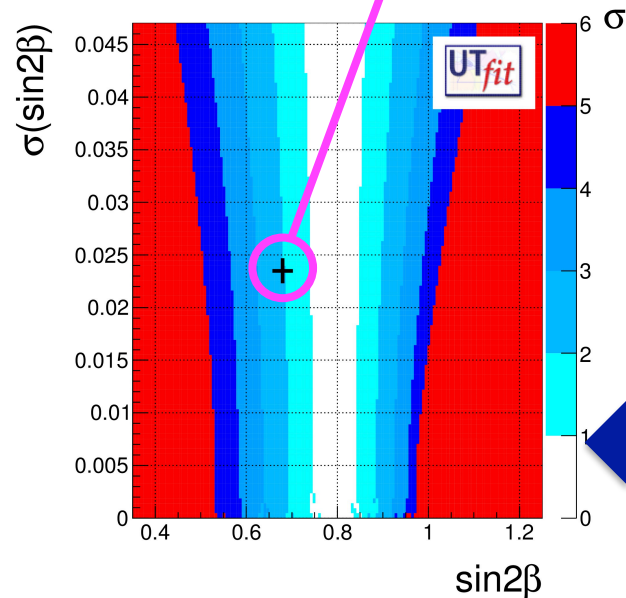
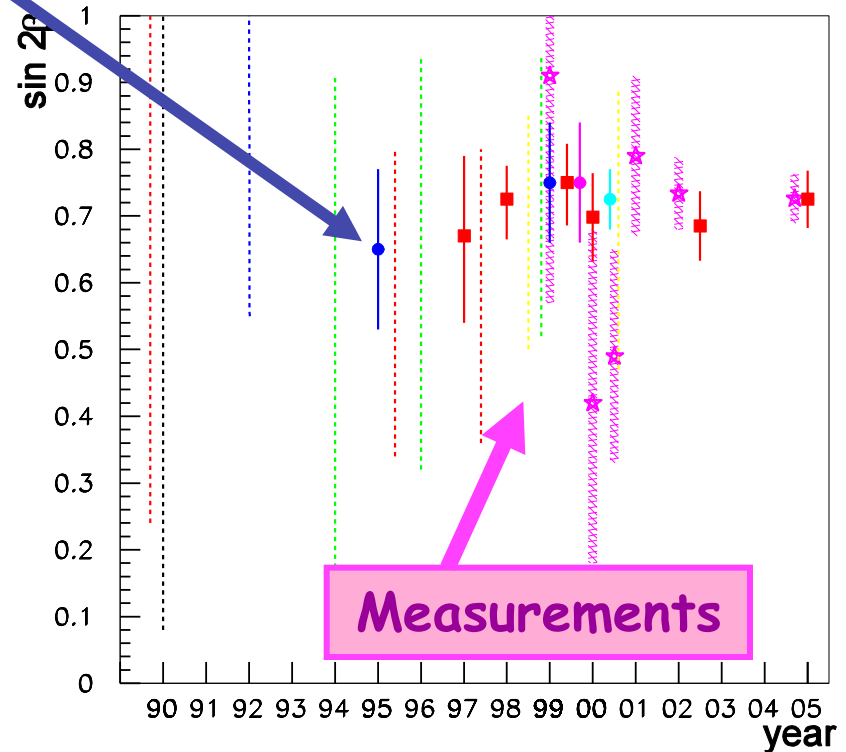
PREDICTION OF $\text{Sin}2\beta$

Ciuchini et al., 1995:
 $\text{Sin}2\beta_{\text{UTA}} = 0.65 \pm 0.12$

Ciuchini et al., 2000:
 $\text{Sin}2\beta_{\text{UTA}} = 0.698 \pm 0.066$

Direct measurement today:
 $\text{Sin}2\beta_{J/\psi K_0} = 0.679 \pm 0.024$

Predictions exist since 1995



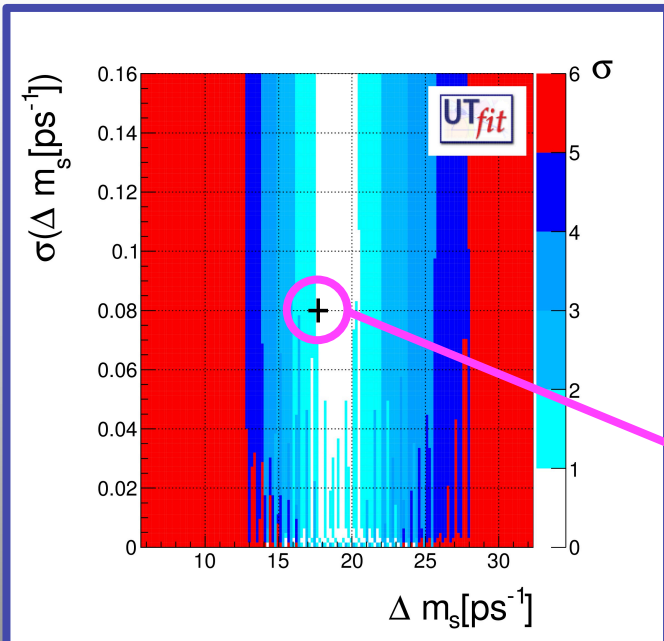
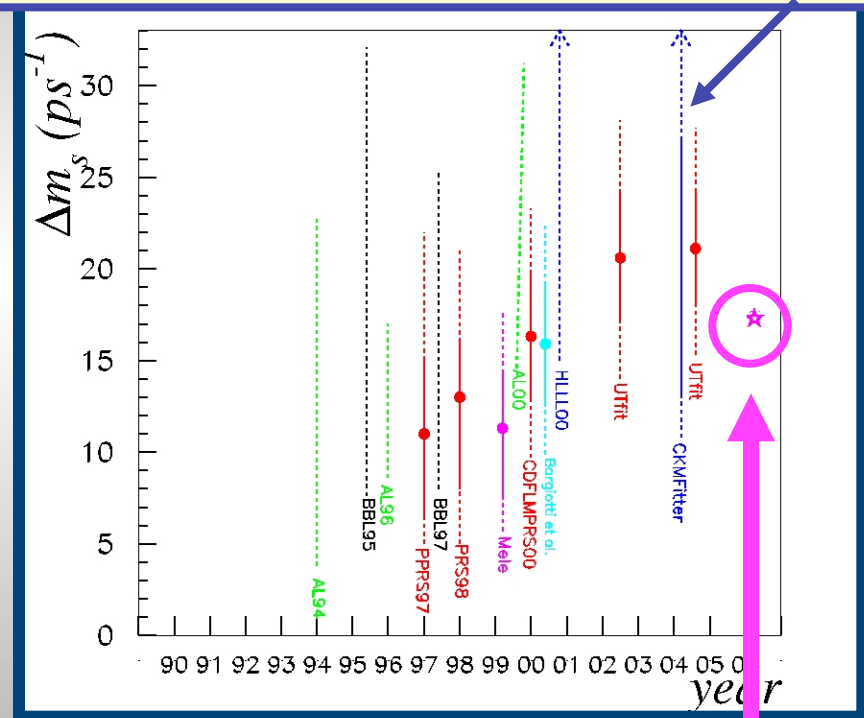
Some tension in the SM fit today:
 $\text{Sin}2\beta_{\text{UTA}} = 0.800 \pm 0.050$

SM PREDICTION OF Δm_s LOOKING FOR NEW PHYSICS EFFECTS

Ciuchini et al., 2000:
 $\Delta m_s = (16.3 \pm 3.4) \text{ ps}^{-1}$

UTfit today:
 $\Delta m_s = (19.0 \pm 1.5) \text{ ps}^{-1}$

The predicted range was very large in the frequentistic CKMFitter approach

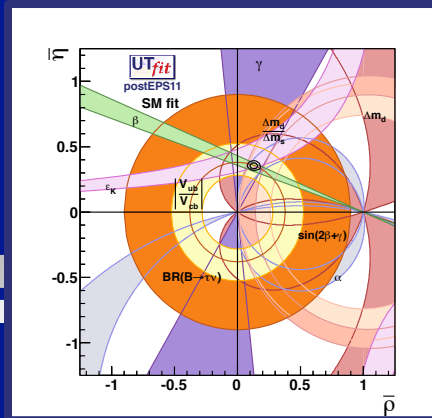


Direct measurement today
 $\Delta m_s = (17.7 \pm 0.08) \text{ ps}^{-1}$

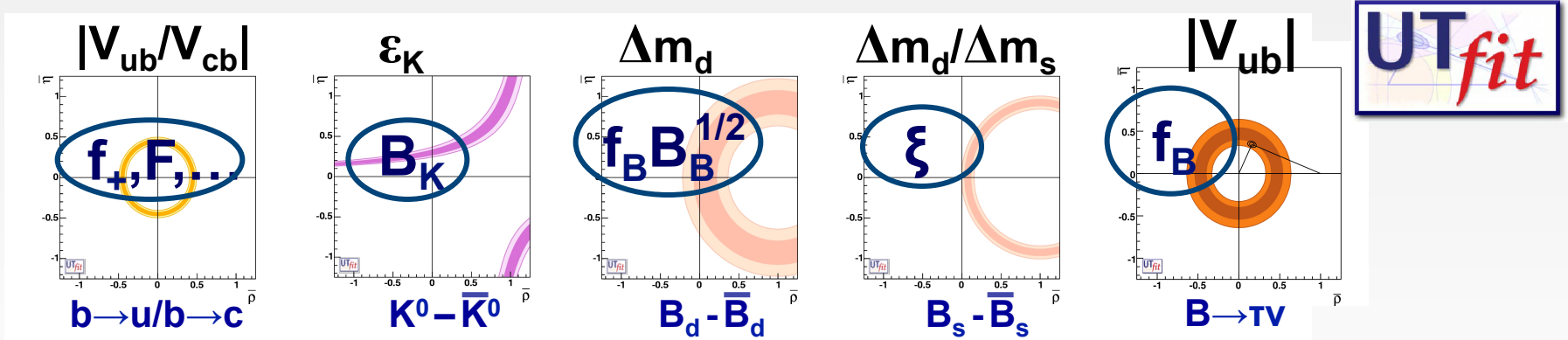
THE UNITARITY TRIANGLE ANALYSIS

② The present

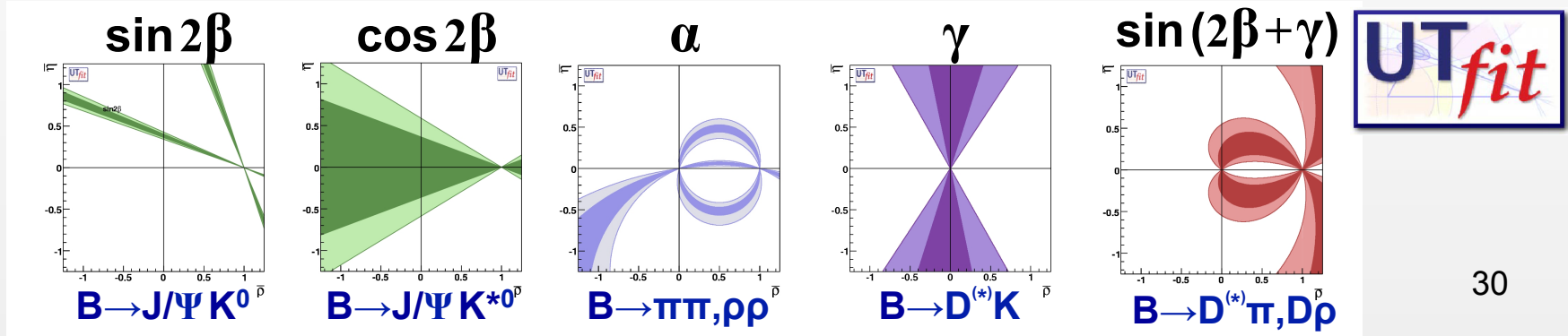
LATTICE QCD AND THE UNITARITY TRIANGLE ANALYSIS



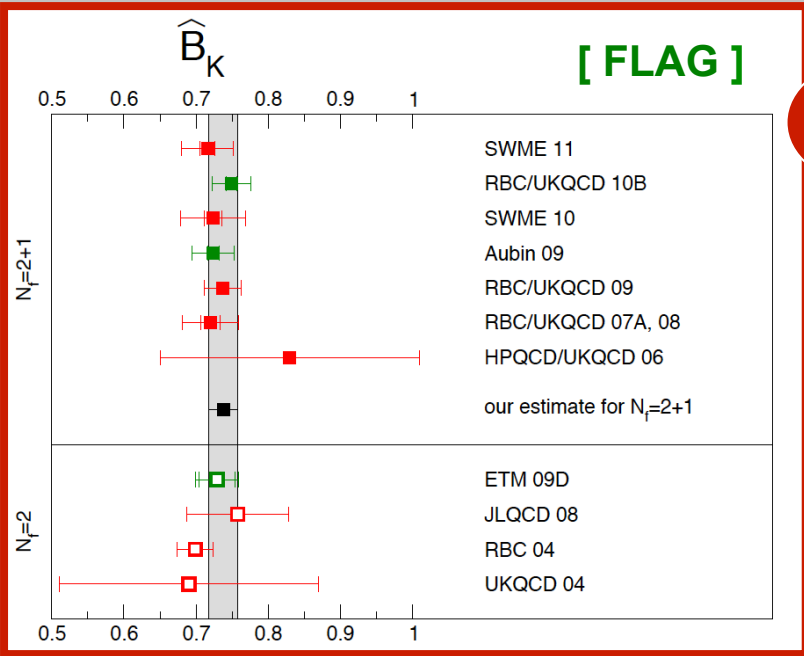
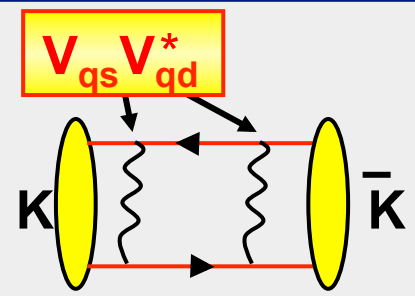
UT-LATTICE



UT-ANGLES



$K^0 - \bar{K}^0$ mixing: B_K



3%

$\hat{B}_K = 0.738(20)$ ($N_f=2+1$)
 $\hat{B}_K = 0.729(30)$ ($N_f=2$) **[FLAG]**

After the FLAG report, a new result has been presented, which claims a remarkable accuracy:

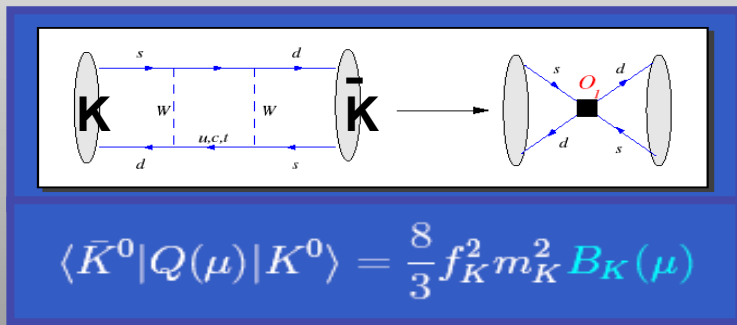
$\hat{B}_K = 0.773(12)$ **BMW' 11**

The average quoted at the latest Lattice conference is:

1.6% !!

$\hat{B}_K = 0.755(12)$

R. Mawhinney @ Lattice 2011

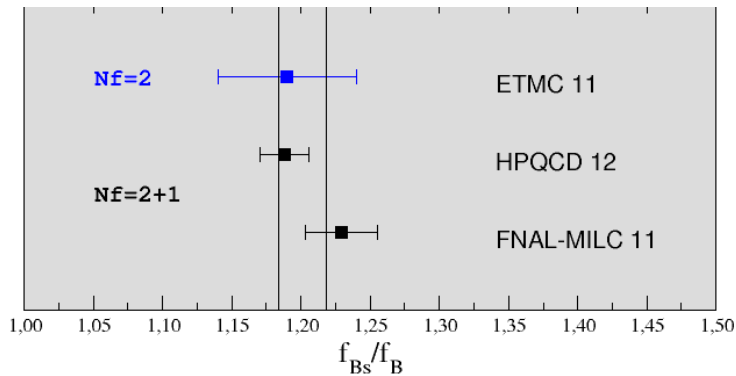
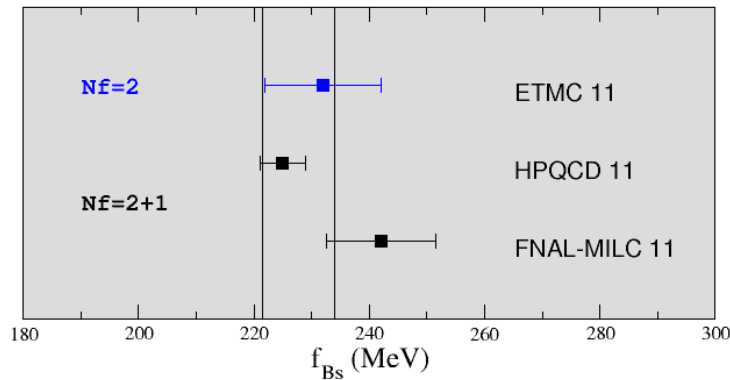
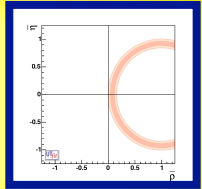


The input currently used in the **UT_{fit}** is:

3%

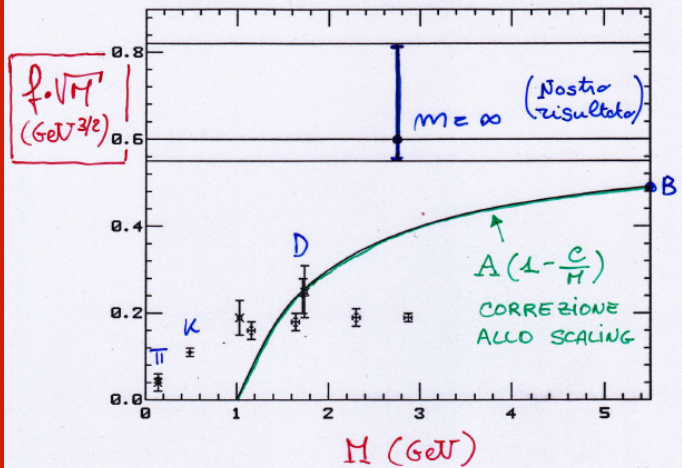
$\hat{B}_K = 0.750(20)$

B-meson decay constants: f_B, f_{B_s}



CORTONA 1990

RISULTATO



- NEL LIMITE ASINTOTICO $M = \infty$:

$$f_B \sqrt{M} \approx (0.60 \pm 0.05 + 0.22) \text{ GeV}^{3/2}$$

$$\rightarrow f_B = (0.26 \pm 0.02 + 0.10) \text{ GeV}$$

- CORREZIONE AL COMPORTAMENTO DI SCALING:

da f_D : $c \approx 1.1 \text{ GeV}$

$$f_B \sqrt{M} = A \left(1 - \frac{c}{M}\right) \rightarrow f_B \approx 200 \text{ MeV}$$

↑ nostro calcolo

2% $f_{B_s} = 227(4) \text{ MeV}$ $f_{B_s} / f_B = 1.192(16)$

C.Davies @ Lattice 2011

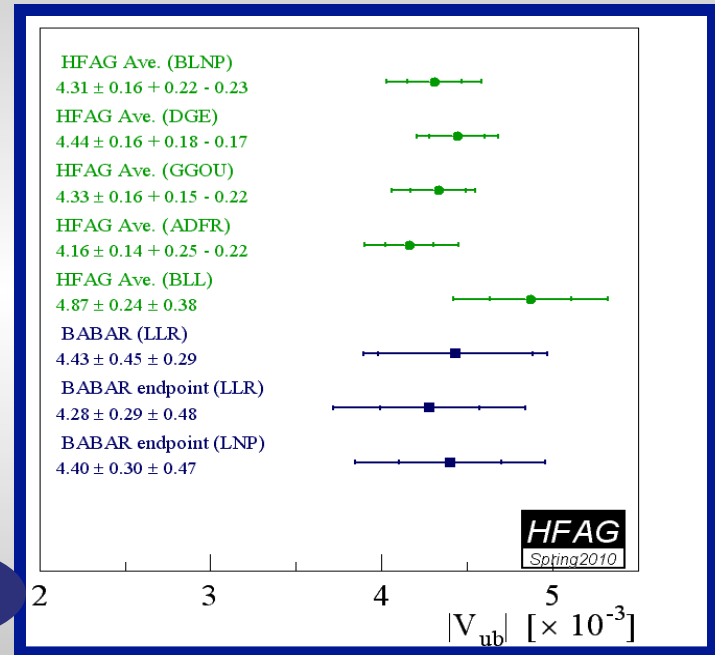
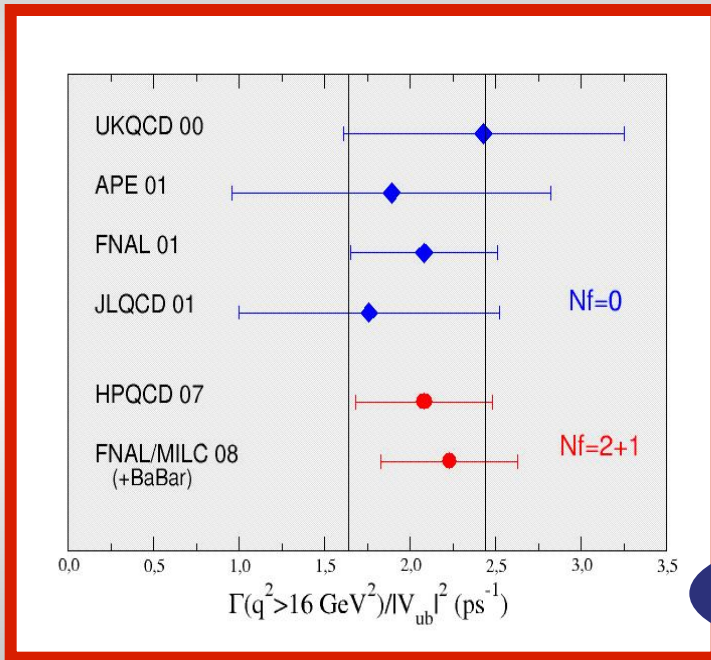
4% $f_{B_s} = 233(10) \text{ MeV}$ $f_{B_s} / f_B = 1.200(20)$

$\rightarrow f_B = 194(9) \text{ MeV}$ used by

Exclusive and Inclusive Vub

THEORETICALLY CLEAN
BUT ONLY TWO MODERN
LATTICE CALCULATIONS

IMPORTANT LONG DISTANCE
CONTRIBUTIONS. THE RESULTS
ARE MODEL DEPENDENT



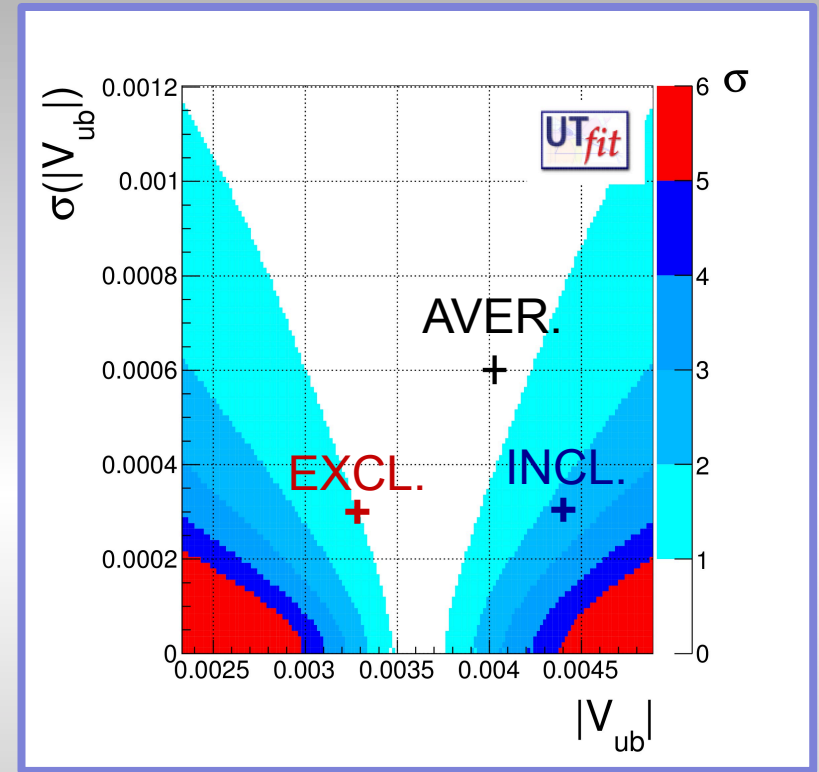
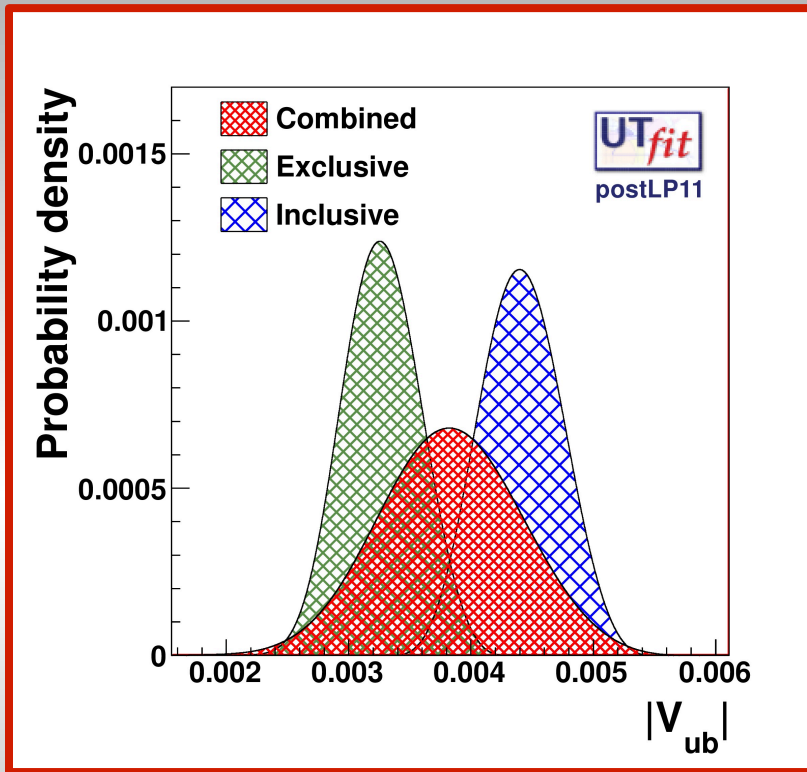
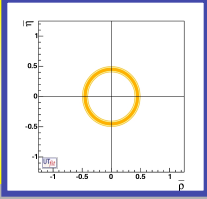
2.6σ

$$|V_{ub}|_{\text{excl.}} = (32.8 \pm 3.0) 10^{-4}$$

$$|V_{ub}|_{\text{incl.}} = (44.0 \pm 3.1) 10^{-4}$$

The uncertainty of inclusive Vub estimated from the spread among different models. This is questionable

Exclusive vs Inclusive V_{ub}

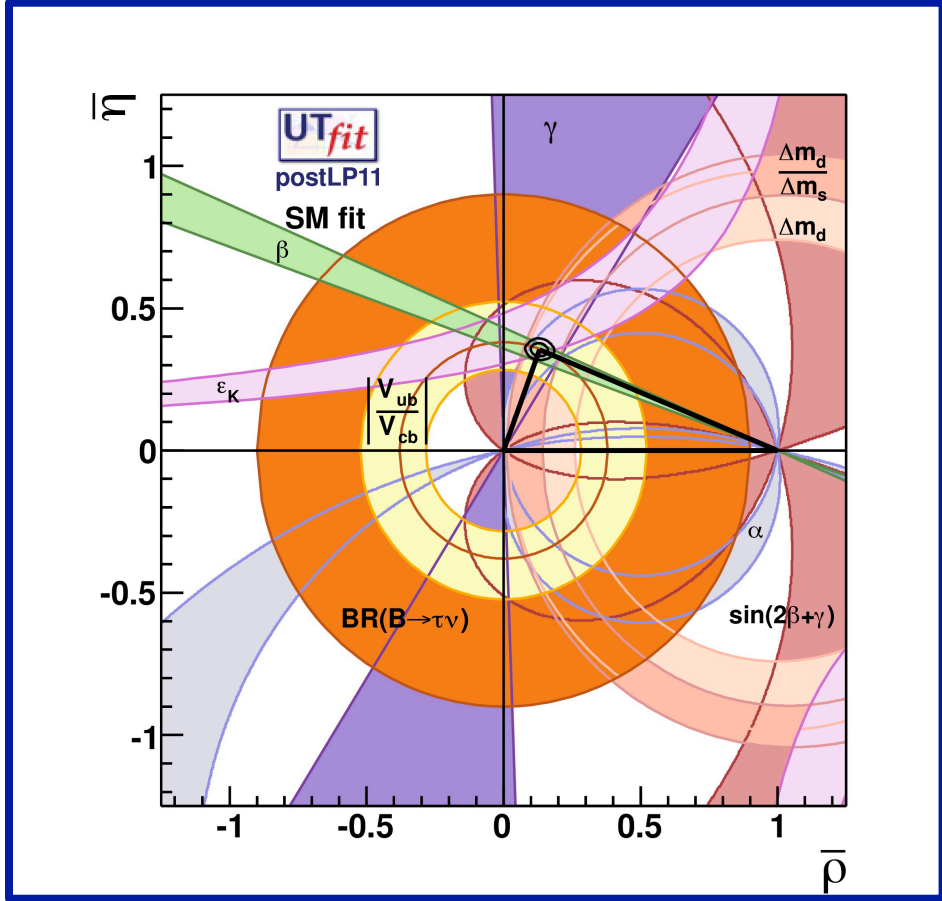


$$|V_{ub}|_{\text{input}} = (38.2 \pm 5.6) 10^{-4}$$

$$|V_{ub}|_{\text{SM-Fit}} = (36.1 \pm 1.4) 10^{-4}$$

Improve the accuracy of exclusive V_{ub} in order to clarify the issue

STANDARD MODEL PREDICTIONS



The CKM Wolfenstein parameters:

$$\bar{\rho} = 0.131 \pm 0.022 \quad 17\%$$

$$\bar{\eta} = 0.354 \pm 0.015 \quad 4\%$$

$$A = 0.817 \pm 0.015$$

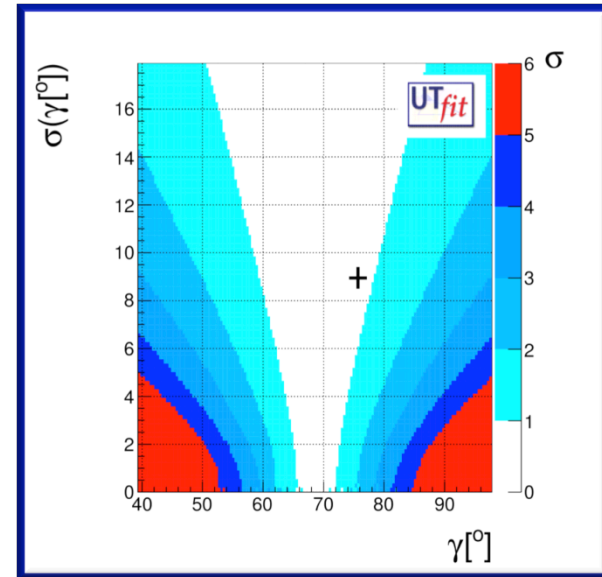
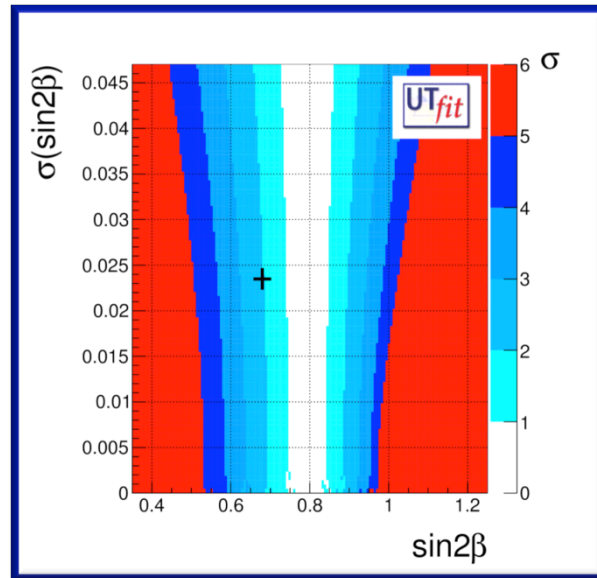
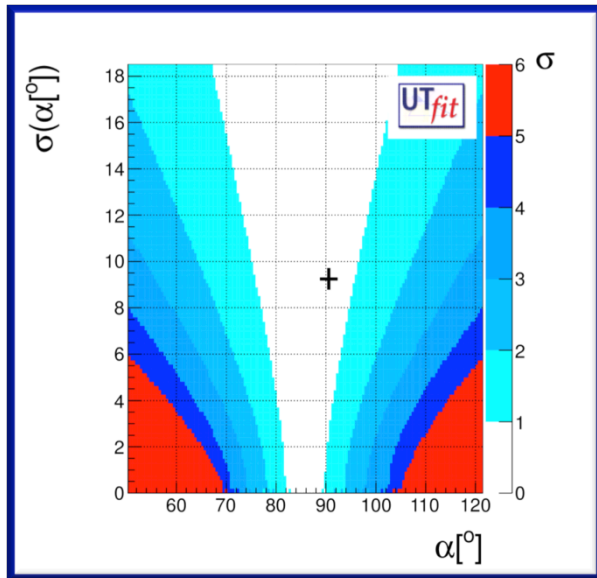
$$\lambda = 0.2252 \pm 0.0010$$

Input value:
 $\lambda = 0.2250 \pm 0.0023$

The fit results for all the nine CKM elements are

$$V = \begin{pmatrix} (0.97427 \pm 0.00012) & (0.22545 \pm 0.00059) & (0.00362 \pm 0.00014)e^{i(-70.0 \pm 3.1)^\circ} \\ (-0.22525 \pm 0.00059)e^{i(0.0349 \pm 0.0015)^\circ} & (0.97338 \pm 0.00012) & (0.0415 \pm 0.00072) \\ (0.00881 \pm 0.00025)e^{i(-22.13 \pm 0.8)^\circ} & (-0.04072 \pm 0.0007)e^{i(1.075 \pm 0.044)^\circ} & (0.999136 \pm 0.00002) \end{pmatrix}$$

The angles:



$(91.4 \pm 6.1)^\circ$
 $(85.8 \pm 3.9)^\circ$

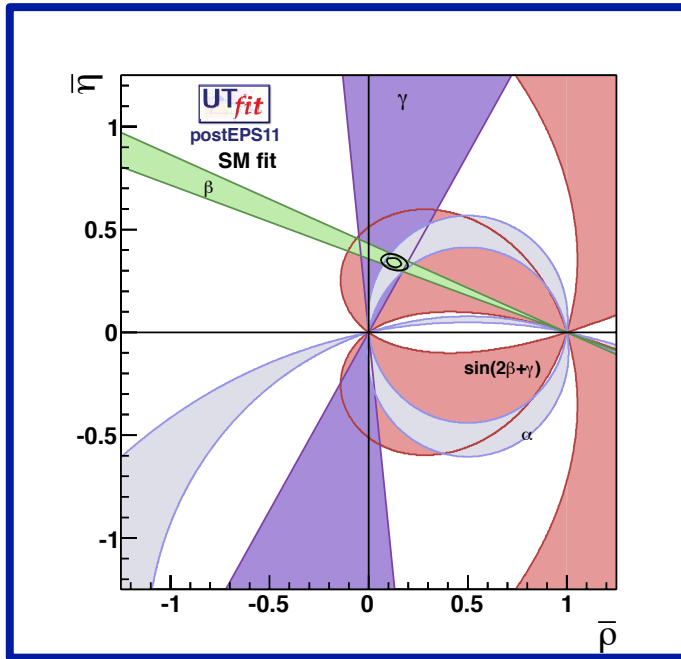
0.679 ± 0.024
 0.80 ± 0.05

2.2 σ tension

$(75.7 \pm 9.2)^\circ$ &
 $(-103.9 \pm 9.2)^\circ$
 $(68.5 \pm 3.2)^\circ$

- Fit input
- SM fit prediction

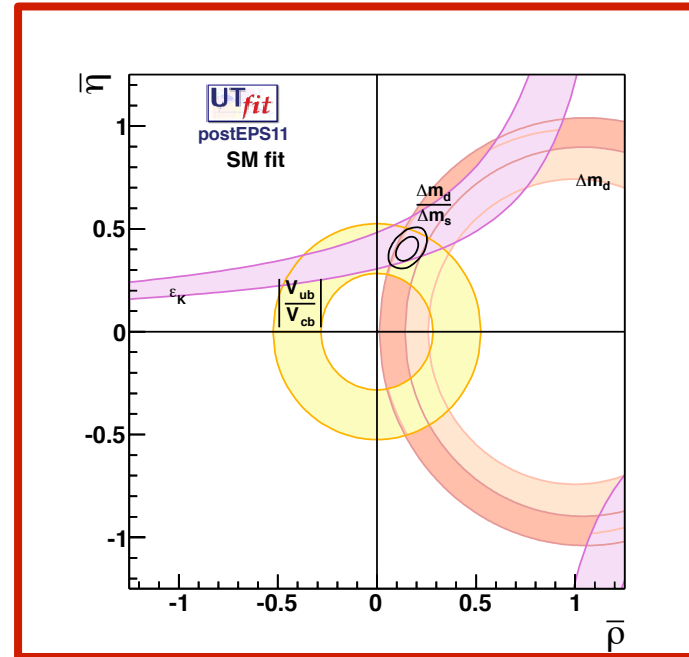
UT-angles



$$\bar{\rho} = 0.129 \pm 0.027$$

$$\bar{\eta} = 0.340 \pm 0.016$$

UT-lattice

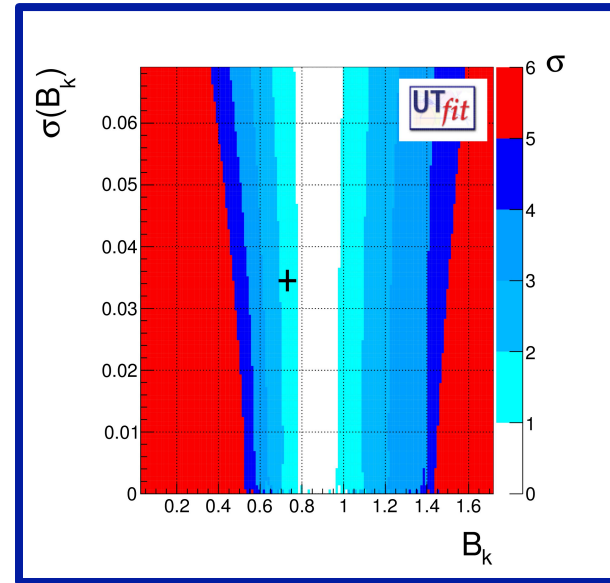
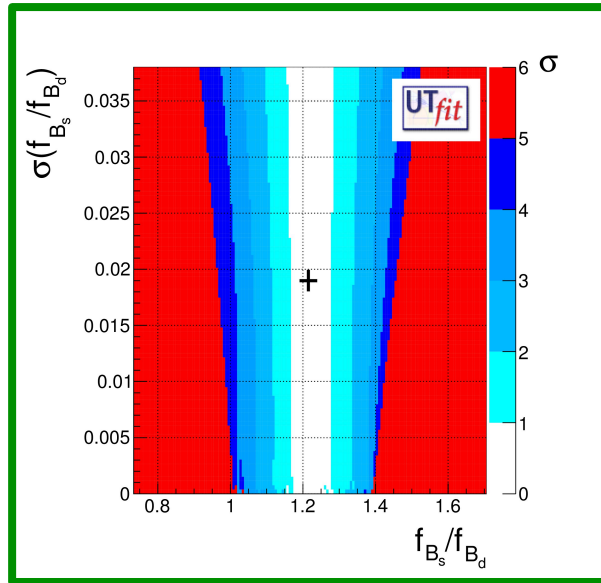
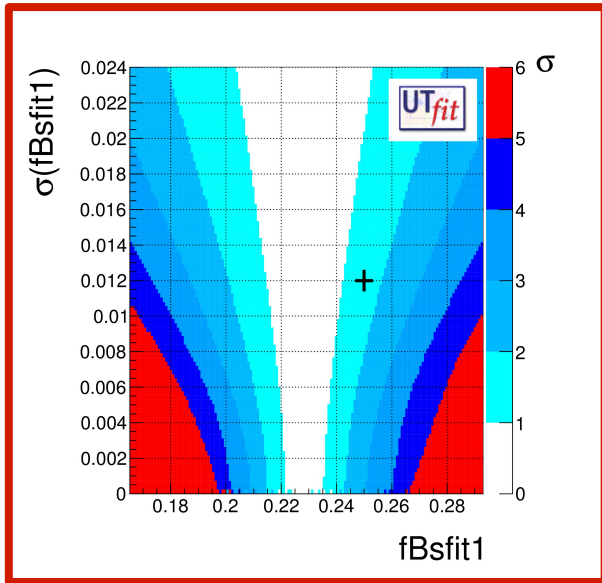


$$\bar{\rho} = 0.155 \pm 0.038$$

$$\bar{\eta} = 0.404 \pm 0.039$$

Lattice inputs are not so relevant today for the Standard Model analysis.
But they are crucial when looking for new physics signals

Assuming the validity of the Standard Model one can perform a fit of the hadronic parameters:



$$f_{B_s} = 250 \pm 12 \text{ MeV [2011]}$$

$$f_{B_s} = 233 \pm 10 \text{ MeV [2012]}$$

$$f_{B_s} = 229 \pm 7 \text{ MeV}$$

$$f_{B_s} / f_{B_d} = 1.215 \pm 0.019 \text{ [2011]}$$

$$f_{B_s} / f_{B_d} = 1.200 \pm 0.020 \text{ [2012]}$$

$$f_{B_s} / f_{B_d} = 1.219 \pm 0.054$$

$$\hat{B}_K = 0.731 \pm 0.035 \text{ [2011]}$$

$$\hat{B}_K = 0.750 \pm 0.020 \text{ [2012]}$$

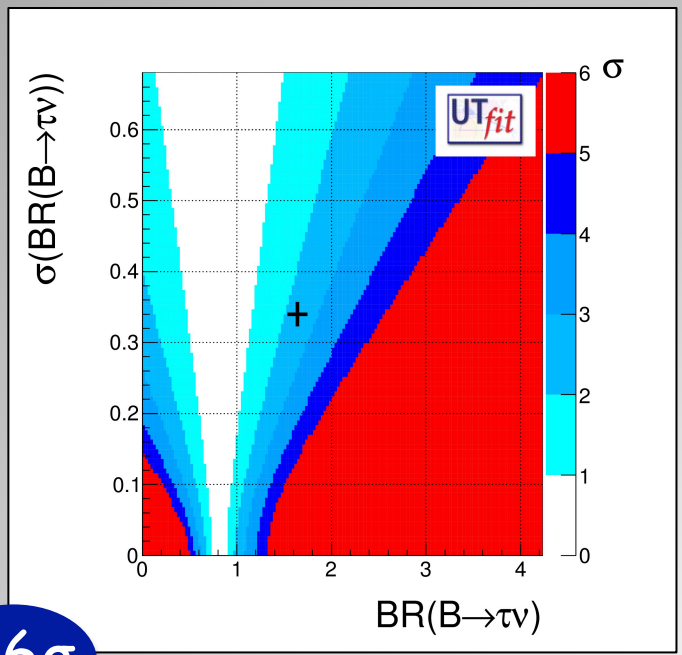
$$\hat{B}_K = 0.872 \pm 0.094$$

- Fit input
- SM fit prediction



Current average:

$$BR(B \rightarrow \tau \nu) = (1.64 \pm 0.34) \times 10^{-4}$$



2.6σ

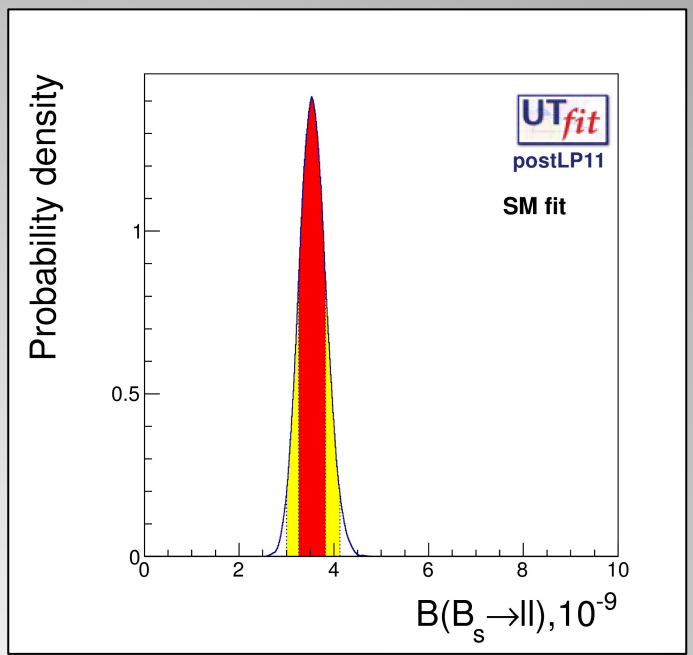
Standard Model determinations from UTfit:

$$BR(B \rightarrow \tau \nu) = (0.83 \pm 0.09) \times 10^{-4}$$



Best upper bound:

$$BR(B_s \rightarrow \mu \mu) < 4.5 \times 10^{-9} @ 95\%CL$$



$$BR(B_s \rightarrow \mu \mu) = (3.54 \pm 0.28) \times 10^{-9}$$

Last minute news from BaBar !!

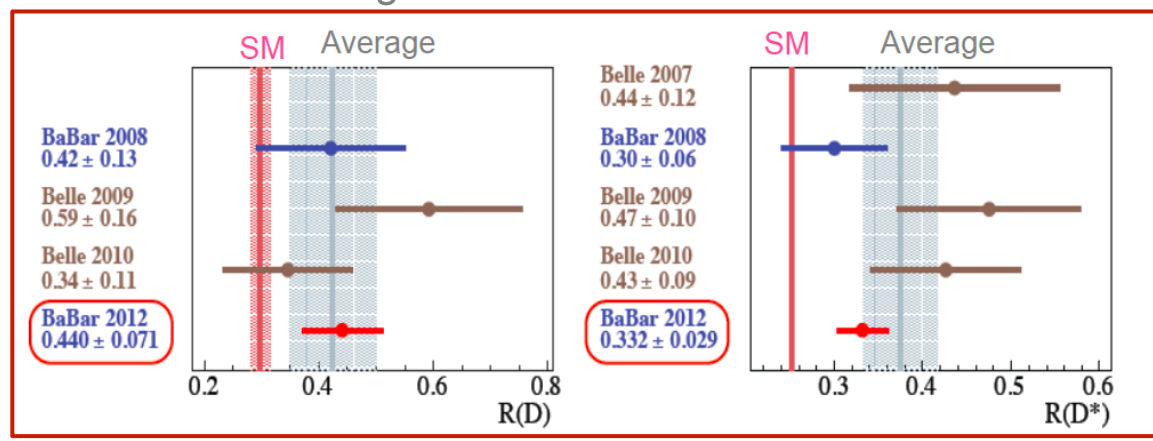
Vera Lüth
SLAC - Stanford University

Presentation at FPCP 2012 @ Hefei (China)
May 21 - 25, 2012

arXiv:1205.5442 [hep-ex] 24 May 2012

$$R(D) = \frac{\Gamma(\bar{B} \rightarrow D\tau\nu)}{\Gamma(\bar{B} \rightarrow D\ell\nu)} \quad R(D^*) = \frac{\Gamma(\bar{B} \rightarrow D^*\tau\nu)}{\Gamma(\bar{B} \rightarrow D^*\ell\nu)}$$

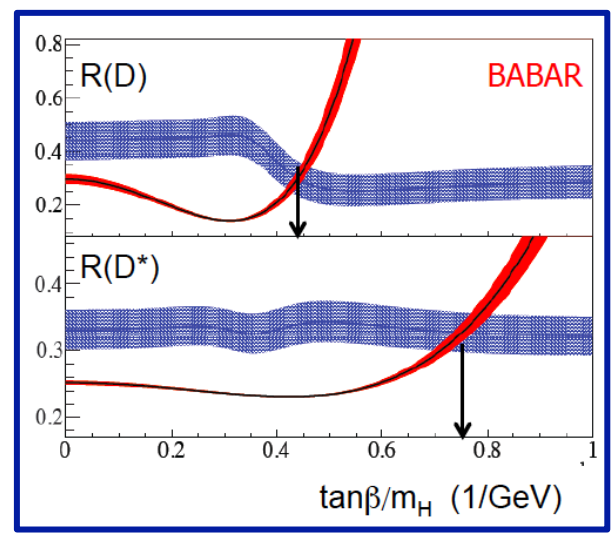
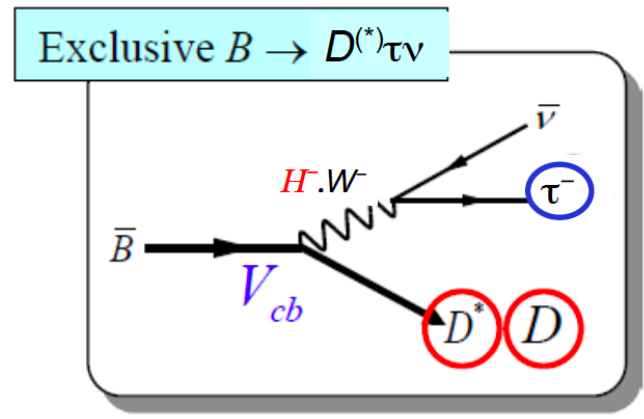
N.B. Average does not include this measurement



The SM prediction is excluded at 3.4σ

$$H_t^{2\text{HDM}} = H_t^{\text{SM}} \times \left(1 - \frac{\tan^2\beta}{m_{H^\pm}^2} \frac{q^2}{1 \mp m_c/m_b} \right)$$

- for $D\tau\nu$
+ for $D^*\tau\nu$



Model-independent determination of the CKM parameters

Assumptions:

- * three generations
- * no NP in tree-level decays
- (* no large NP contributions to EW penguin in $B \rightarrow \pi\pi$ and Γ_q)

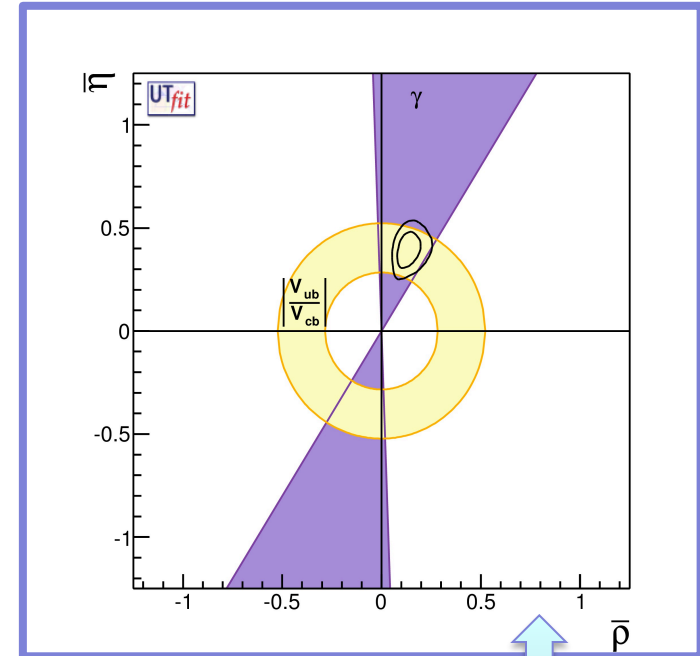
$$\bar{\rho} = 0.134 \pm 0.044$$

$$\bar{\eta} = 0.403 \pm 0.058$$

In the SM was

$$\bar{\rho} = 0.131 \pm 0.022$$

$$\bar{\eta} = 0.354 \pm 0.015$$



The degeneracy on γ broken by A_{SL}

A number of additional constraints are included: semileptonic asymmetries (A_{SL}^d, A_{SL}^s), lifetime differences and mixing phases ($\Delta\Gamma_d/\Gamma_d, \Delta\Gamma_s/\Gamma_s, \varphi_s$).

NP in mixing amplitudes parameterized in a general form:

- K mixing amplitude
(2 real parameters)

$$\text{Re } A_K = C_{\Delta m_K} \text{Re } A_K^{SM}$$

$$\text{Im } A_K = C_{\varepsilon_K} \text{Im } A_K^{SM}$$

- B_d and B_s mixing amplitudes
(2+2 real parameters)

$$A_q e^{2i\phi_q} = C_{Bq} e^{2i\phi_{Bq}} A_q^{SM} e^{2i\phi_q^{SM}}$$

Observables:

In the Standard Model: $C_{xx} = 1$, $\phi_{xx} = 0$

$$\Delta m_{q/K} = C_{Bq/\Delta m_K} (\Delta m_{q/K})^{SM}$$

$$a_{CP}^{B_d \rightarrow J/\psi K_s} \rightarrow \sin 2(\beta + \phi_{B_d})$$

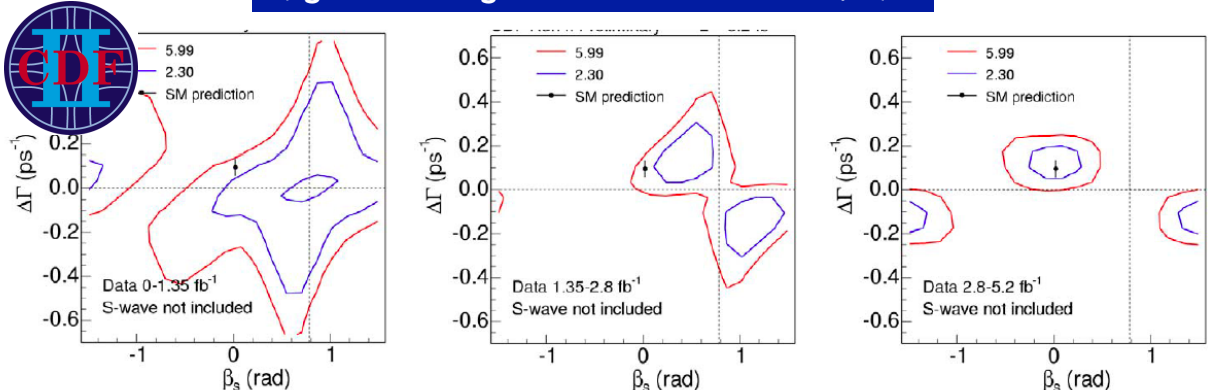
$$a_{SL}^q = \text{Im}(\Gamma_{12}^q / A_q)$$

$$\varepsilon_K = C_{\varepsilon} \varepsilon_K^{SM}$$

$$a_{CP}^{B_s \rightarrow J/\psi \phi} \rightarrow -\beta_s + \phi_{B_s}$$

$$\Delta \Gamma^q / \Delta m_q = \text{Re}(\Gamma_{12}^q / A_q)$$

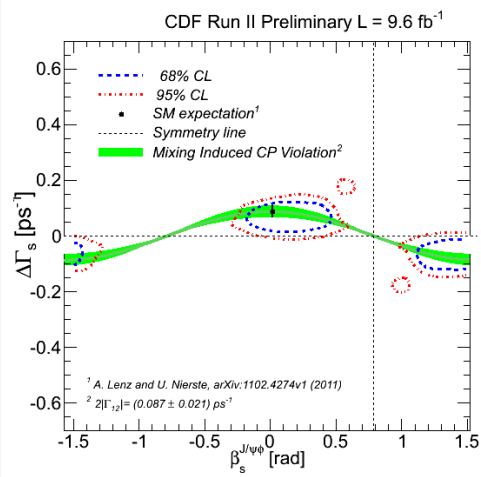
ϕ_s vs $\Delta\Gamma_s$ from $B_s \rightarrow J/\psi \phi$



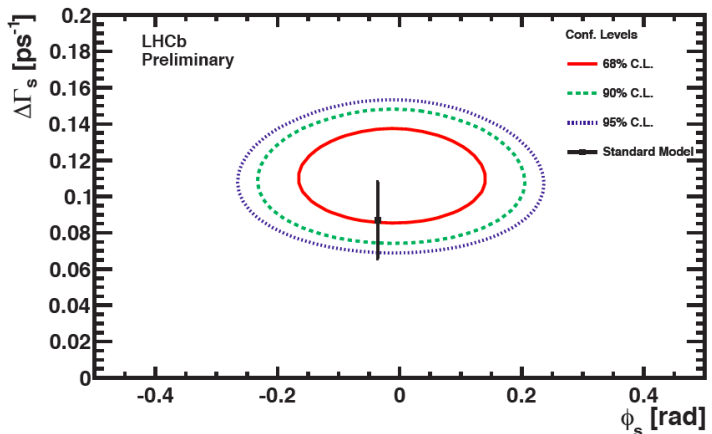
CDF: end of 2007

Added in 2008

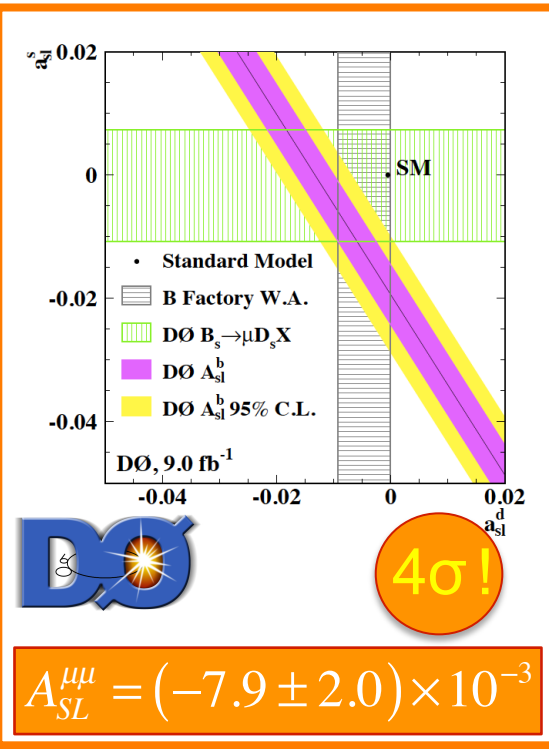
Added in 2010



Full data sample

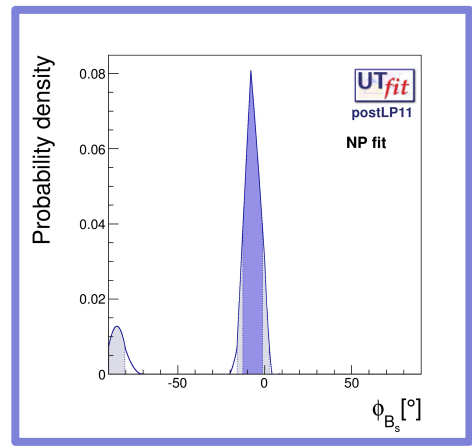


LHCb

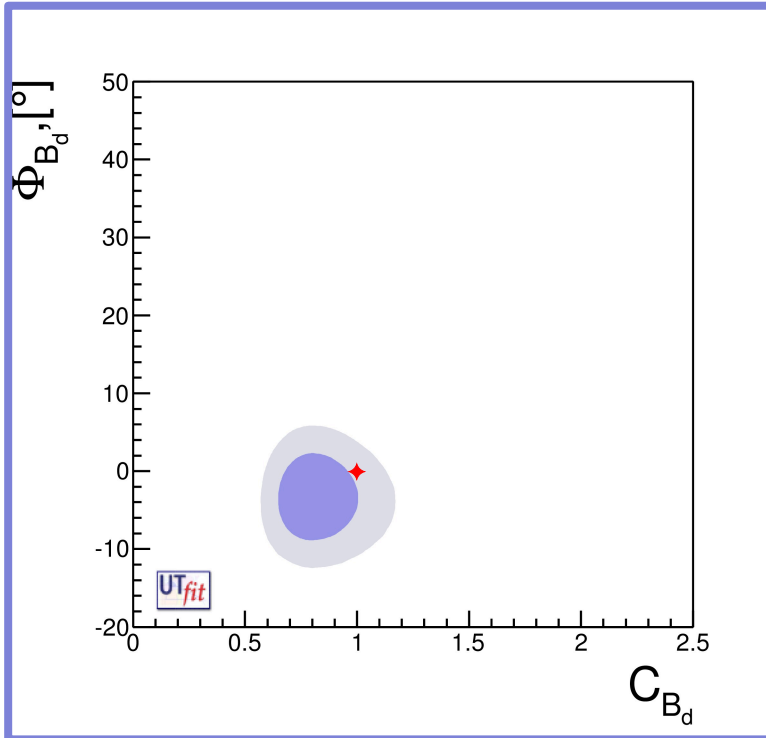


4 σ !

$$A_{SL}^{\mu\mu} = (-7.9 \pm 2.0) \times 10^{-3}$$



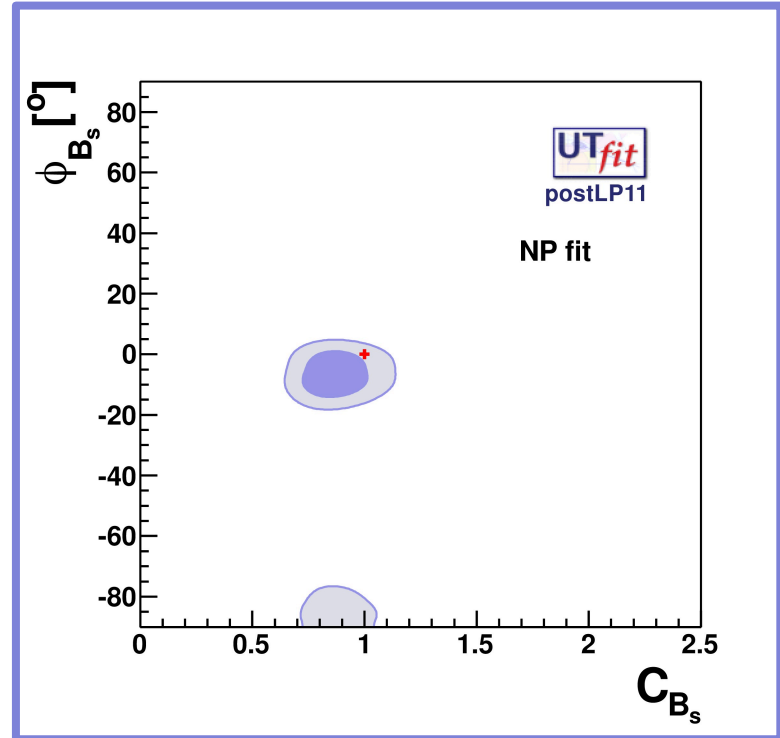
B_d - B_d mixing



$$C_{B_d} = 0.81 \pm 0.12$$

$$\varphi_{B_d} = (-3.4 \pm 3.7)^\circ$$

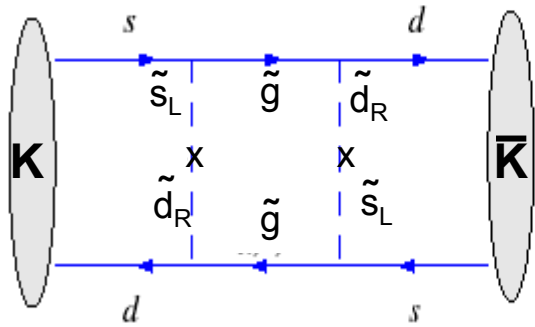
B_s - B_s mixing



$$C_{B_s} = 0.87 \pm 0.10$$

$$\varphi_{B_s} = (-6.9 \pm 5.6)^\circ$$

K-K̄ MIXING BSM

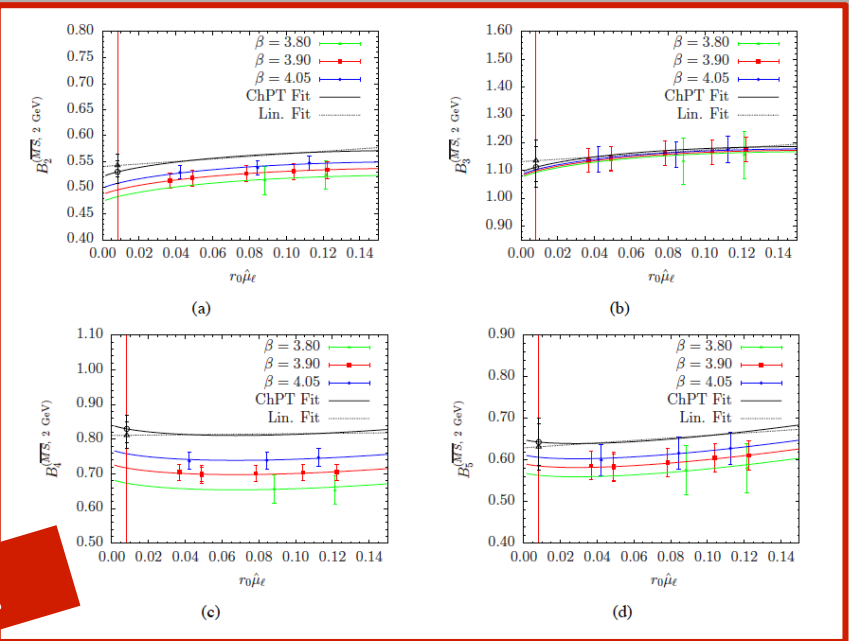


$$\begin{aligned} \mathcal{O}_1 &= [\bar{s}^\alpha \gamma_\mu (1 - \gamma_5) d^\alpha] [\bar{s}^\beta \gamma_\mu (1 - \gamma_5) d^\beta] \\ \mathcal{O}_2 &= [\bar{s}^\alpha (1 - \gamma_5) d^\alpha] [\bar{s}^\beta (1 - \gamma_5) d^\beta] \\ \mathcal{O}_3 &= [\bar{s}^\alpha (1 - \gamma_5) d^\beta] [\bar{s}^\beta (1 - \gamma_5) d^\alpha] \\ \mathcal{O}_4 &= [\bar{s}^\alpha (1 - \gamma_5) d^\alpha] [\bar{s}^\beta (1 + \gamma_5) d^\beta] \\ \mathcal{O}_5 &= [\bar{s}^\alpha (1 - \gamma_5) d^\beta] [\bar{s}^\beta (1 + \gamma_5) d^\alpha] \end{aligned}$$



in preparation

Previous results in the quenched approximation only (and quite in disagreement)



B_i (RI-MOM at 2 GeV)

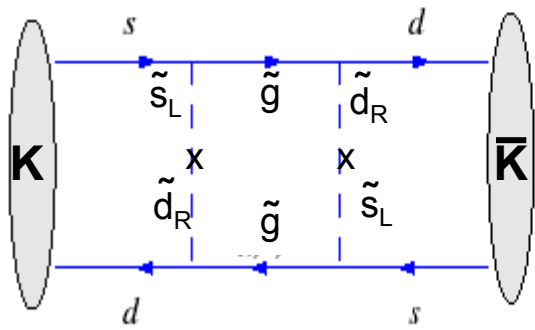
		Ref. [6]		Ref. [5]	
This work		$a = 0.09$ fm	$a = 0.13$ fm	$a = 0.07$ fm	$a = 0.09$ fm
CL					
1	0.51(02)	0.56(05)	0.53(04)	0.68(21)	0.70(15)
2	0.73(04)	0.87(07)	0.90(10)	0.67(07)	0.72(09)
3	1.29(11)	1.41(12)	1.53(40)	0.95(15)	1.21(10)
4	1.04(07)	0.94(05)	0.90(13)	1.00(09)	1.15(05)
5	0.76(09)	0.62(05)	0.56(14)	0.66(11)	0.88(06)

ETMC 2012

Babich et al. 2006

Donini et al. 1999

K-K̄ MIXING BSM



LOWER BOUNDS ON THE NP SCALE:

UTfit 0707.0636

$$C_i(\Lambda) = F_i L_i / \Lambda^2$$

L_i = loop factor
 F_i = flavour coupling
 (e.g. $L_i = (\alpha_s)^2$ for gluino exchange in the MSSM)

$$\langle K | H_{\text{eff}}^{\text{NP}} | \bar{K} \rangle = \sum_{i,j} C_j(\Lambda) W_{ji}(\Lambda, \mu) \underbrace{\langle K | Q_i(\mu) | \bar{K} \rangle}_{\text{LATTICE QCD}}$$

NP SCALE \nearrow

$$L_i = F_i = 1$$

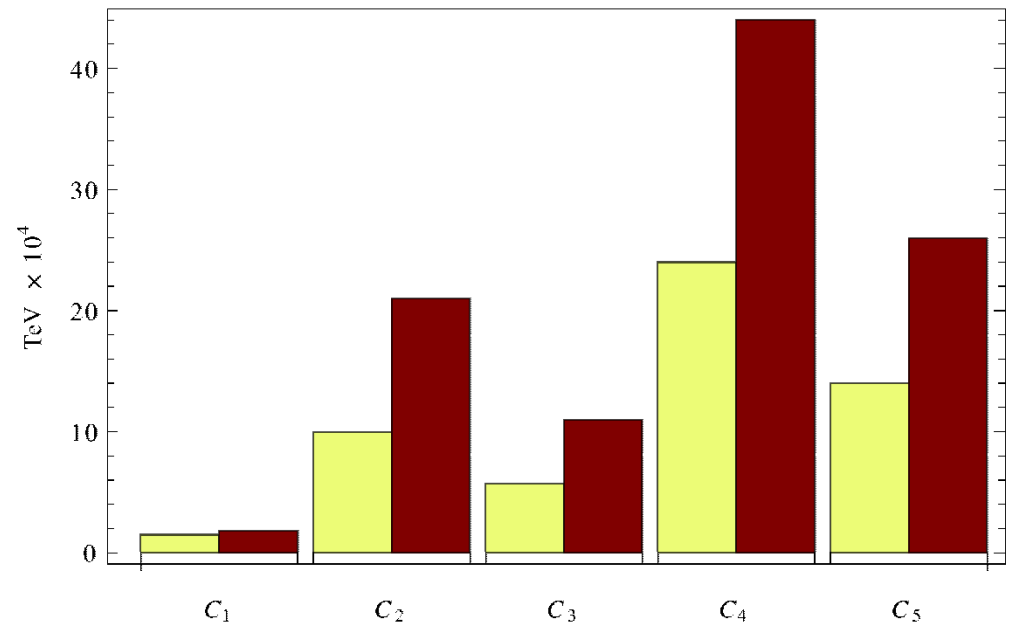
	95% allowed range (GeV ⁻²)	Lower limit on Λ (TeV)
$\text{Im } C_1^K$	$[-2.7, 3.0] \cdot 10^{-15}$	$1.8 \cdot 10^4$
$\text{Im } C_2^K$	$[-2.3, 2.2] \cdot 10^{-17}$	$21 \cdot 10^4$
$\text{Im } C_3^K$	$[-8.0, 8.4] \cdot 10^{-17}$	$11 \cdot 10^4$
$\text{Im } C_4^K$	$[-5.0, 5.1] \cdot 10^{-18}$	$44 \cdot 10^4$
$\text{Im } C_5^K$	$[-1.5, 1.5] \cdot 10^{-17}$	$26 \cdot 10^4$

$$\alpha_s(\Lambda) \Lambda \sim 4 \cdot 10^4 \text{ TeV}$$

$$\alpha_W \Lambda \sim 10^4 \text{ TeV}$$

$$\sqrt{F_{\text{SM}}} \Lambda \sim 10^2 \text{ TeV}$$

$$\sqrt{F_{\text{SM}}} \alpha_W \Lambda \sim 3 \text{ TeV}$$



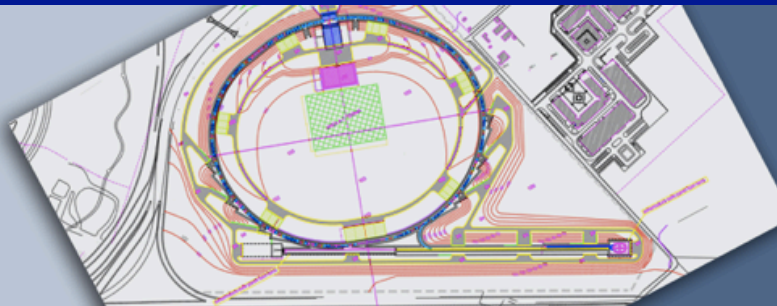
THE UNITARITY TRIANGLE ANALYSIS

③ The future
The SuperB era

La SuperB e il Cabibbo-Lab



A A A



- Home
- Scienza
- Progetto
- Il Laboratorio
- Strumenti
- Ricadute tecnologiche
- Stampa e Media
- Lavorare al Cabibbo Lab
- Collaborazione SuperB

News

- [Il CabibboLab apre le porte a nuovi fisici, ingegneri e tecnici](#)
- [Un super laser per SuperB](#)
- [SuperB: Scelto il team che costruirà l'acceleratore](#)

cerca...

MEETINGS

31/5/2012 - 5/6/2012
4th SuperB Collaboration Meeting -
Isola d'Elba Italy

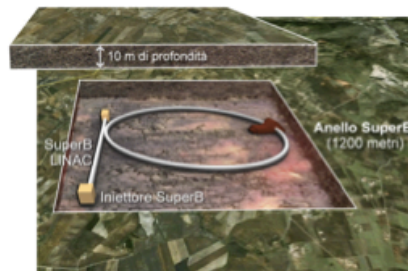
19/3/2012 - 23/3/2012
3rd SuperB Collaboration Meeting -
LNF

13/12/2011 - 16/12/2011
2nd SuperB Collaboration Meeting -
LNF

Il Cabibbo-Lab



Il Cabibbo-Lab è il centro di ricerca internazionale per la **fisica fondamentale e applicata** che occuperà un'area di circa 30 ettari nel campus dell'Università di Roma Tor Vergata. Proposto dall'Istituto Nazionale di Fisica Nucleare (INFN), il progetto condurrà entro cinque anni alla costruzione dell'acceleratore SuperB, uno dei più significativi tra i 14 **progetti bandiera** del Piano di Ricerca Nazionale del MIUR, approvato dal CIPE.

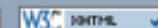


Prima macchina acceleratrice ideata per soddisfare allo stesso tempo le esigenze della fisica fondamentale e di quella applicata, **SuperB** è il cuore del Cabibbo-Lab: un tunnel sotterraneo di 1,3 km di circonferenza dove si scontreranno elettroni e positroni, con l'obiettivo di fare luce su alcuni dei più affascinanti **interrogativi di fisica contemporanea** e rendere disponibili allo stesso tempo sorgenti di luce con caratteristiche tali da permettere di studiare la materia organica e le nanostrutture.

In perfetta sinergia con l'acceleratore LHC del Cern di Ginevra, il progetto SuperB si appresta ad affrontare una nuova frontiera della fisica sperimentale delle alte energie,

aumentando a livelli senza precedenti l'intensità delle collisioni tra le particelle accelerate e garantendo in questo modo la produzione di **fenomeni fisici rarissimi** e ancora inesplorati. SuperB permetterà così lo studio dei meccanismi che hanno prodotto la scomparsa dell'**antimateria** poco dopo il Big Bang e le forze che tengono uniti i componenti fondamentali della materia.

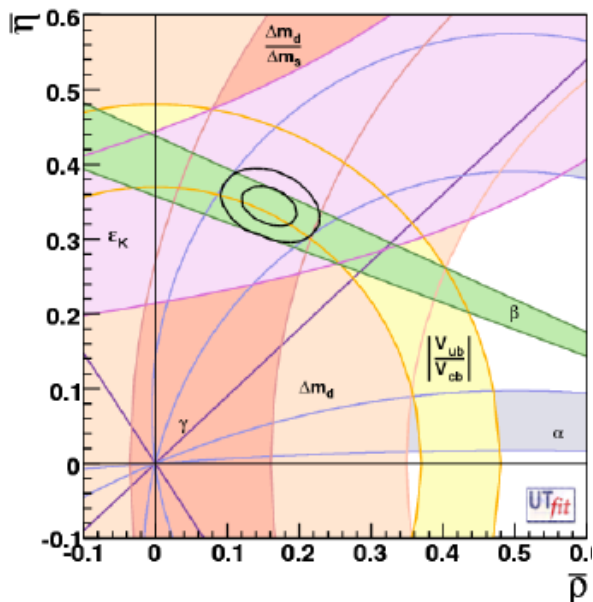
A lato degli obiettivi di fisica fondamentale, da subito un'ampia comunità scientifica interdisciplinare, italiana e internazionale beneficerà della possibilità di utilizzare la **luce di sincrotrone** emessa dagli elettroni nell'acceleratore. Si tratta di fasci di luce con caratteristiche uniche per coerenza e collimazione, tali da consentire di visualizzare strutture biologiche o inorganiche a una risoluzione mai raggiunta, e di scattare delle **"microistantanee"** dei processi biochimici in atto. Offrendo una storica



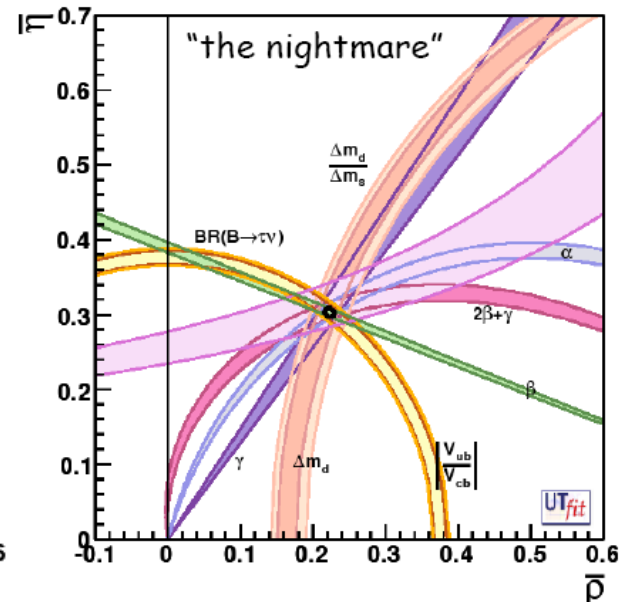
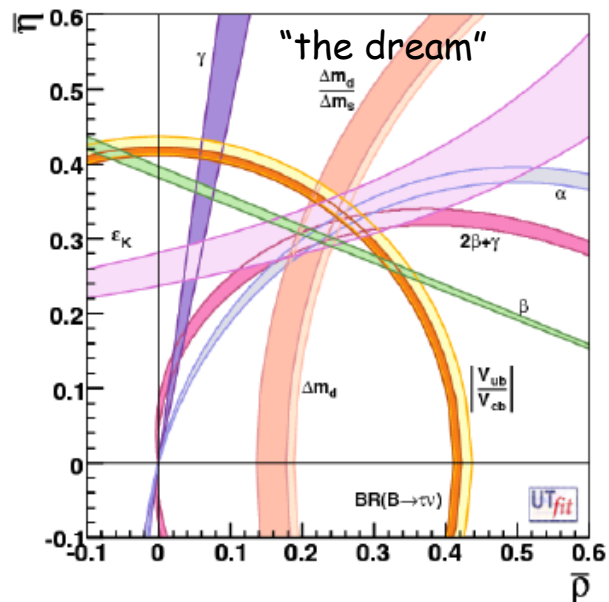
LQCD at a SuperB factory (2016-)

For example: testing the CKM paradigm at the 1% level

Today



With a SuperB in 2016



The theoretical accuracy must compete with the experimental one.

Can we reach the 1% accuracy in Lattice QCD ??

Cost of the "SuperB" lattice simulation

Simulation parameters

Nconf = 120

$a = 0.033$ fm
[$1/a = 6.0$ GeV]

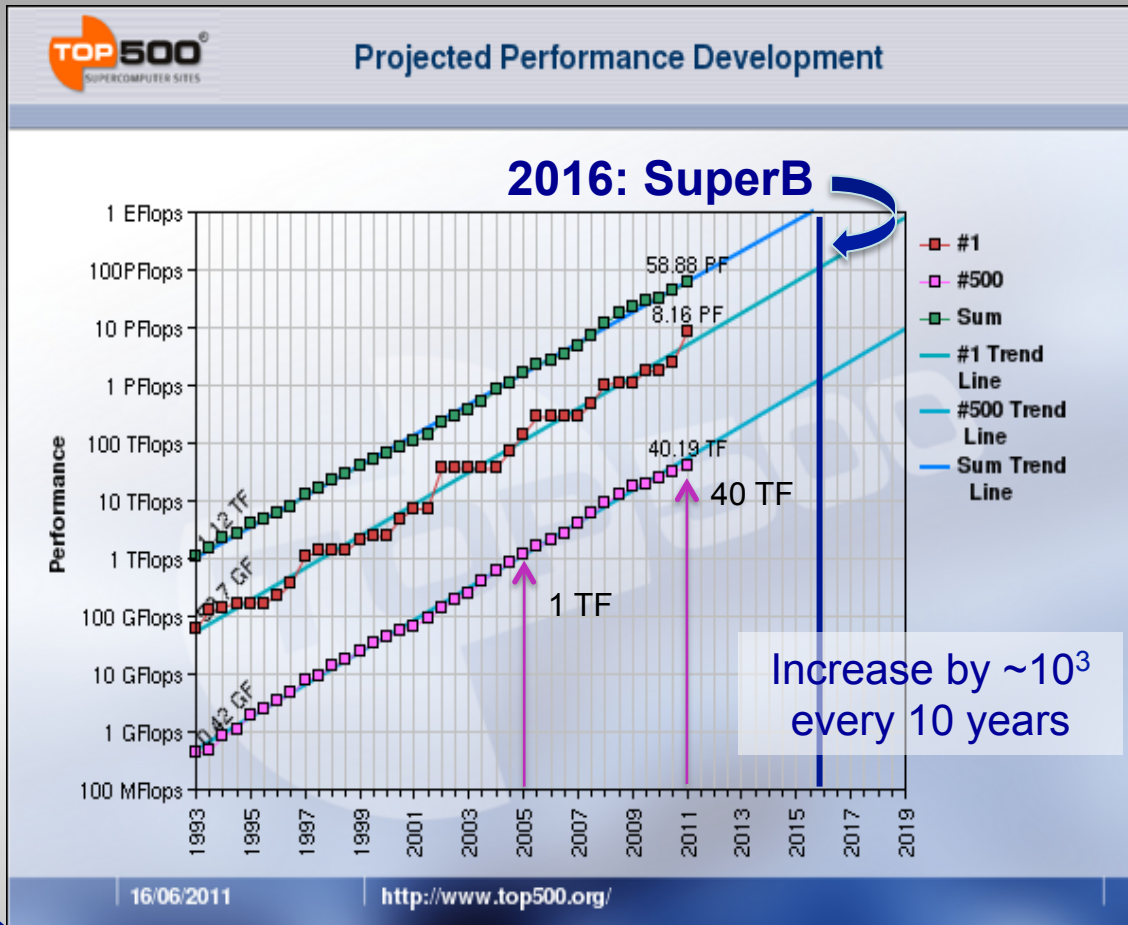
$\hat{m}/m_s = 1/12$

[$M_\pi = 200$ MeV]

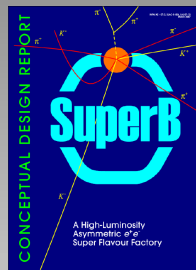
$L_s = 4.5$ fm
[$V = 136^3 \times 270$]

~ 3 PFlop-years

VL @ SuperB IV



Affordable with
1-10 PFlops available for
Lattice QCD in 2016 !

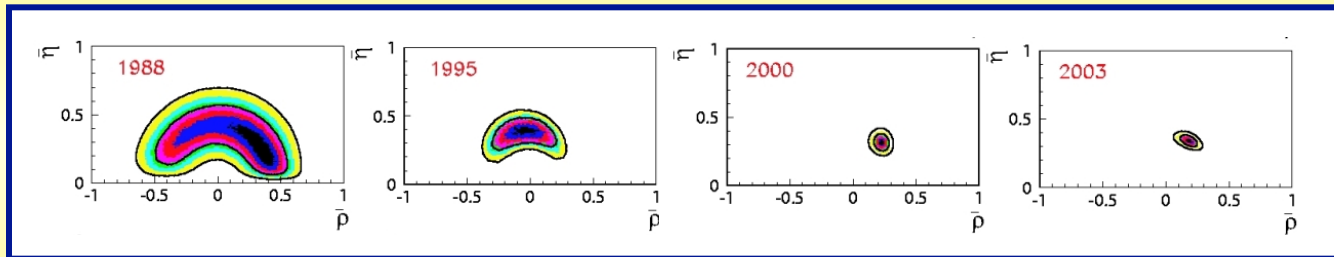


VL, SuperB CDR,
arXiv:0709.0451
updated in arXiv:1008.1541

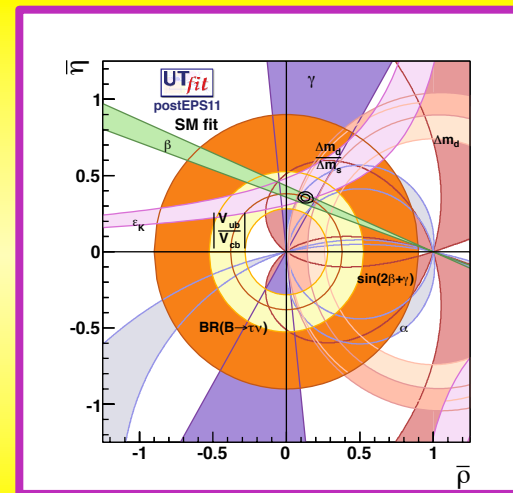
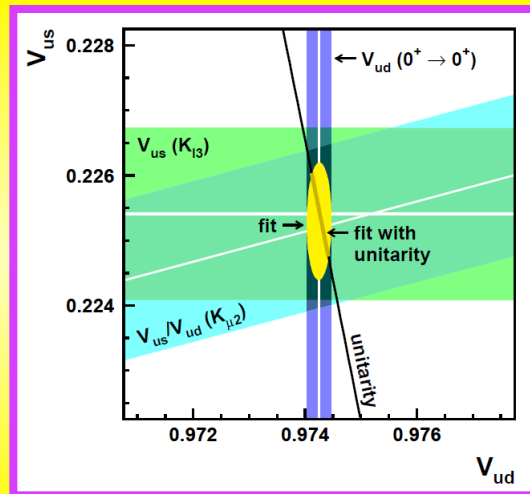


Hadronic matrix element	Lattice error in 2006	Lattice error in 2009	6 TFlop Year [2009]	60 TFlop Year [2012 LHCb]	1-10 PFlop Year [2016 SuperB]
$f_+^{K\pi}(0)$	0.9%	0.5%	0.7%	0.4%	< 0.1%
\hat{B}_K	11%	5%	5%	3%	1%
f_B	14%	5%	3.5 - 4.5%	2.5 - 4.0%	1 - 1.5%
$f_{B_s} B_{B_s}^{1/2}$	13%	5%	4 - 5%	3 - 4%	1 - 1.5%
ξ	5%	2%	3%	1.5 - 2 %	0.5 - 0.8 %
$\mathcal{F}_{B \rightarrow D/D^*lv}$	4%	2%	2%	1.2%	0.5%
$f_+^{B\pi}, \dots$	11%	11%	5.5 - 6.5%	4 - 5%	2 - 3%
$T_1^{B \rightarrow K^*/\rho}$	13%	13%	----	----	3 - 4%

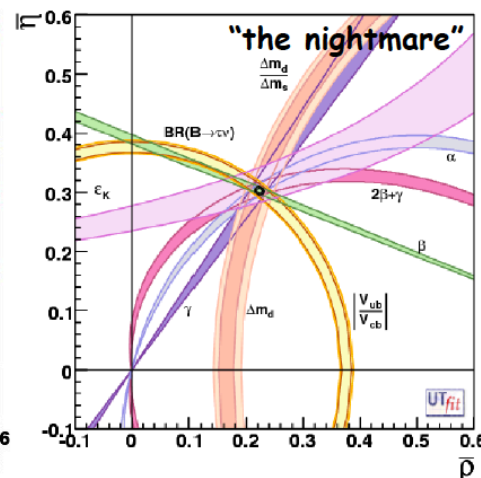
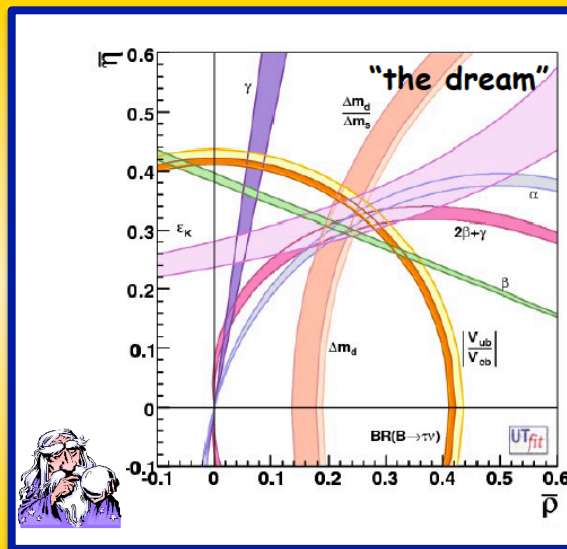
The past



the present



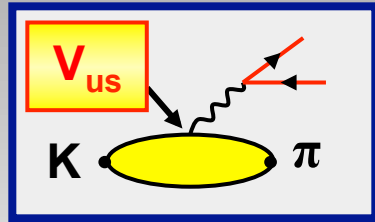
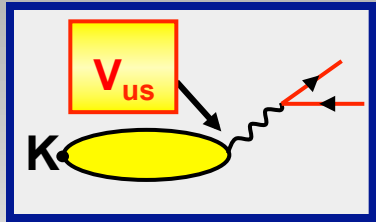
and the future



2016 @ SuperB

Supplementary slides

LQCD independent estimates of f_K/f_π , $f_+(0)$



$$\frac{|V_{us}|}{|V_{ud}|} \frac{f_K}{f_\pi} = 0.2758(5) \quad 1$$

$$|V_{us}| f_+(0) = 0.2163(5) \quad 2$$



Assuming the Standard Model and combining with nuclear β decays

$$|V_{ud}|^2 + |V_{us}|^2 + \cancel{|V_{ub}|^2} = 1 \quad 3$$

$$|V_{ud}| = 0.97425(22) \quad 4$$

From 20 superallowed transitions
[Hardy and Towner 08]

one obtains:

$$f_K/f_\pi = 1.192(6)$$

$$f_+(0) = 0.960(5)$$

and:

$$|V_{us}| = 0.2254(10)$$

The error is at the per mille level: a challenge for Lattice QCD

FLAG [Flavianet Lattice Averaging Group] (1st report)

G.Colangelo, S.Dürr, A.Jüttner, L.Lellouch, H.Leutwyler, V.Lubicz, S.Necco,
C.Sachrajda, S.Simula, T.Vladikas, U.Wenger, H.Wittig

Exclusive vs Inclusive Vcb

EXCLUSIVE: 2 APPROACHES

- "double ratios" (FNAL)
- "step scaling" (TOV)

Good agreement

$$|V_{cb}|_{\text{excl.}} = (39.5 \pm 1.0) 10^{-3}$$

$$|V_{cb}|_{\text{incl.}} = (41.7 \pm 0.7) 10^{-3}$$

1.8 σ

